

Abrupt climate and weather changes across time scales

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Key Points:

- Non-linear response of the climate system to external orbital forcing
- Centennial and millennial variability are stochastic in nature
- Challenges of ocean-radiocarbon modeling with high-resolution at the coasts and high latitudes
- The past provides evidences of abrupt climate changes and weather extremes

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Abstract

The past provides evidence of abrupt climate shifts and changes in the frequency of climate and weather extremes. We explore the non-linear response to orbital forcing and then consider climate millennial variability down to daily weather events. Orbital changes are translated into regional responses in temperature, where the precessional response is related to nonlinearities and seasonal biases in the system. We question regularities found in climate events by analyzing the distribution of inter-event waiting times. Periodicities of about 900 and 1150 years are found in ice cores besides the prominent 1500-years cycle. However, the variability remains indistinguishable from a random process, suggesting that centennial-to-millennial variability is stochastic in nature. New numerical techniques are developed allowing for a high resolution in the dynamically relevant regions like coasts, major upwelling regions, and high latitudes. Using this model, we find a strong sensitivity of the Atlantic meridional overturning circulation depending on where the deglacial meltwater is injected into. Meltwater into the Mississippi and near Labrador hardly affect the large-scale ocean circulation, whereas subpolar hosing mimicking icebergs yields a quasi shutdown. The same multi-scale approach is applied to radiocarbon simulations enabling a dynamical interpretation of marine sediment cores. Finally, abrupt climate events also have counterparts in the recent climate records, revealing a close link between climate variability, the statistics of North Atlantic weather patterns, and extreme events.

Plain Language Summary

Predicting the future spread of possible climates, the risk of climate extremes and the risk of rapid transitions is of high socio-economic relevance. The past provides evidence of abrupt climate change and the frequency of extremes. This allows to separate anthropogenic signals from natural climate variability. Earth system models applied both to past and future scenarios will enhance our ability to detect regime shifts which are necessary to potentially predict climate extremes and transitions. We consider the response of the system to regular orbital forcing and then focus on shorter time scales down to weather. The appearance of precession is linked to non-linear responses of the climate system to external orbital forcing. Furthermore, we find that centennial-to-millennial variability is stochastic in nature. We also discuss recent developments of climate models with superior resolution in typical retrieval regions of paleoclimate records, such as continental margins and coasts. Using this model, we find a strong sensitivity

38 of the Atlantic meridional overturning circulation depending on where the deglacial meltwater is injected
39 into. Meltwater into the Mississippi and near Labrador hardly affect the large-scale ocean circulation,
40 whereas subpolar hosing related to icebergs yields a quasi shutdown. Our multi-scale approach is applied
41 to radiocarbon simulations enabling a dynamical interpretation of marine sediment cores. Finally, we ex-
42 plore a close link between climate variability, the statistics of weather patterns, and extreme events.

1 Introduction

Weather and climate vary on broad ranges of spatial and temporal scales. This is evident from observations and simulations of the present climate as well as from climate history as recorded in geological and glaciological archives (Mitchell, 1976; Crowley & North, 1991; Saltzman, 2002; Bradley, 2014). Climate and Earth system models are widely used to evaluate the impact of anthropogenic emissions on future climate. The validation of these models by simulating different climate scenarios is essential to understand the sensitivity of the climate system to external forcing. The models are clearly unrivaled in their ability to simulate a broad range of large-scale phenomena on seasonal to decadal time scales (Flato et al., 2013). However, the reliability of models to simulate climate variability on multi-decadal and longer time scales requires additional evaluation (Haywood et al., 2019). Climate records derived from paleoenvironmental parameters facilitate the testing of models out of the "comfort zone of present-day climate", or, in other words, one application of paleoclimate is to validate state-of-the-art coupled climate models for past time slices and past climate transitions.

In this review, we will highlight key areas where major future efforts are needed to address fundamental questions related to abrupt changes of climate and weather patterns across a range of time scales (Fig. 1). As climate can rapidly change, potential thresholds for the occurrence of climate extremes and transitions must be estimated. During the most recent glacial period, the Northern Hemisphere was subject to abrupt large-scale transitions between cold (stadial) and warmer (interstadial) conditions (Fig. 2a). These shifts are well documented by isotope records from Greenland ice cores, and are known as Dansgaard-Oeschger (DO) events (Dansgaard et al., 1993). They typically show an abrupt warming of up to 10 – 15 K within a few decades followed by a gradual cooling that stretches over several hundred to several thousand years. We evaluate the stochastic nature of this multi-centennial-to-millennial climate variability and will present a Bayesian methodology to quantify the uncertainties. Such concepts are also useful to interpret decadal and multi-decadal variability, linked to convection events in the Labrador Sea.

The pacing of regional to global climate variations on time scales of tens of thousands of years is related to changes in the seasonal and latitudinal sunlight distribution induced by variations in the Earth's orbital parameters (e.g., Milankovitch, 1941; Berger, 1978; Laskar et al., 2004). A key element of the orbital theory is that summer insolation at high latitudes of the Northern Hemisphere determines glacial-interglacial transitions connected with the waxing and waning of large continental ice sheets (e.g. Im-

72 brie and Imbrie, 1980). During the last two million years, these glacial-interglacial cycles provide the dom-
73 inant signal in the climate record. The annual mean insolation forcing shows the obliquity cycle at a time
74 scale of ~ 40 ky (Fig. 2b). While precessional forcing (~ 20 ky variations; ky=1000 years) is clearly vis-
75 ible for seasonal forcing (Fig. 2a), it is zero for local annual mean insolation forcing (Fig. 2b). Therefore,
76 the precessional signal is an indicator for non-linear responses to the forcing (such as stronger response
77 in summer than winter) or for non-local responses, which raises the question of forcing and feedback mech-
78 anisms for paleoclimatic change (Short et al., 1991; Brickmann et al., 1999; Valdes & Glover, 1999; Tziper-
79 man et al., 2006; Laepple & Lohmann, 2009). We are going to reveal the response to orbital forcing in
80 a complex climate model. This model can be viewed as a benchmark test of the climate system, espe-
81 cially for the question of how orbital changes are translated into local temperature variations. This reg-
82 ular, external orbital forcing is strongly linked to changes in seasonality (Fig. 1).

83 Paleoclimatic evidence suggests that some past climate shifts were associated with changes in North
84 Atlantic Deep Water (NADW) formation at the end of the last ice-age (Lehman & Keigwin, 1992; Dans-
85 gaard et al., 1993; Sarnthein et al., 1994). During the last deglaciation about 20-10 ky before present,
86 the climate has warmed due to insolation and rising carbon dioxide (CO_2) concentrations (Fig. 2b,c). Con-
87 siderable meltwater discharge from the large continental ice sheets over the Northern Hemisphere entered
88 the North Atlantic. The most prominent meltwater event occurred at about 14.6 ky BP (BP=before present)
89 and was named Meltwater Pulse 1a (MWP-1a) (Fairbanks, 1989) with up to 0.5 Sv ($1 \text{ Sv} = 10^6 \text{ m}^3 \text{ s}^{-1}$)
90 (Clark et al., 2002a) entering the North Atlantic with some possible contributions from the Arctic Ocean
91 (Tarasov & Peltier, 2006) and Southern Ocean (Carlson & Clark, 2012; Weber et al., 2014). However,
92 paleoclimate data indicate that NADW formation was not affected seriously (McManus et al., 2004), and
93 the next major cooling phases occurred either earlier (during Heinrich-event 1 at about 17.5 ky BP) or
94 more than 1,000 years later at the onset of the Younger Dryas between about 12.9 ky BP and 11.5 ky
95 BP (Clark et al., 2002b). Analyses of coarse-grained ice-rafted debris and planktonic foraminifers revealed
96 pronounced dropstone layers that have been deposited in the North Atlantic documenting ice rafting dur-
97 ing cold events (Heinrich, 1988). At the end of the Younger Dryas, the North Atlantic realm again ex-
98 perience rapid warming that ended up in Holocene climate conditions. In the discussion of probable causes,
99 the melting and calving of continental ice sheets is the obvious source of the sea-level changes (e.g. Flower
100 et al., 2004; Carlson, 2009; Otto-Bliesner & Brady, 2010). Therefore, we assess the effect of freshwater

101 originating from the Northern Hemisphere ice sheets on the ocean circulation. It turns out that the way
102 the freshwater enters the ocean, matters: In high-resolution models, the coastal hosing can be trapped
103 with minor influences on ocean circulation. In contrast, solid freshwater (calving ice bergs) could have
104 been transported off the coasts affecting deep water formation effectively. This sensitivity is essential to
105 understand past and potential future abrupt climate changes related to freshening of the North Atlantic
106 Ocean (e.g., Kjeldsen et al., 2015; Sejr et al., 2017).

107 One important aspect of paleoclimate is that it can avoid overconfidence of our limited view onto
108 the current system. A climatologist without knowledge of paleoclimate records, such as the ice-core curve
109 in Fig. 2a before the last 9,000 years, could assume that the local temperature over Greenland follows
110 the boreal summer insolation. However, when understanding climate dynamics and predicting the future
111 spread of possible climates, we find that the comprehension of irregular centennial-to-millennial climate
112 variations and abrupt shifts is essential. Knowledge regarding the system can be obtained either from
113 proxies recording past climate and environmental conditions, and/or by simulating the climate to get in-
114 formation on mechanisms under various forcings to identify leads and lags in the system.

115 One of the greatest obstacles of resolving the leads and lags-puzzle is the difficulty of developing
116 an accurate time scale for long ice cores and marine sediments. The development of such a time scale would
117 allow for testing of many climate-forcing hypotheses, leading to significant advances in paleoclimate the-
118 ory. In the late Quaternary, the backbone of such a time scale could be provided through radiocarbon
119 (^{14}C) dating. Numerical models are key to quantify processes affecting ^{14}C based age models of climate
120 records. This is particularly the case for marine climate records. While the ocean is the most important
121 reservoir of heat and carbon regarding decadal to millennial-scale climate variability, a direct evaluation
122 of ^{14}C -dated marine records is hampered by various effects introducing dating uncertainties. Prior to the
123 Holocene, these effects are poorly constrained and have to be inferred either from ad-hoc assumptions
124 or through modeling, which may involve ensemble approaches to explore its uncertainty range (e.g. Butzin
125 et al., 2012a, 2017, 2020). To show an example, the blue curve in Fig. 3 shows a conversion from ^{14}C years
126 to calendar years without the information guided by numerical modeling of the ocean-carbon system (e.g.,
127 Reimer et al., 2004). Using a ^{14}C -equipped climate model where paleoceanographic information is assim-
128 ilated (Butzin et al., 2012a), the black curve in Fig. 3 displays a non-stationary relation between the sim-
129 ulated ^{14}C years to calendar years relation. The difference between the resulting dates can be up to 1000

130 years and can vary in time. A further complication in radiocarbon dating is that marine records typi-
131 cally originate from continental margins, marginal seas, or tropical islands. Such sites may not be rep-
132 resentative to mirror large- or global-scale processes of climatic change, demanding a valid extrapolation
133 method to infer past global conditions, which makes a proper data-model comparison difficult (Lohmann
134 et al., 2013a). High-resolution models are required to elucidate the causal chains in the climate system,
135 notably during abrupt transitions of the last deglaciation, and provide a benchmark for future transitions
136 under rapid CO₂ increase (cf. Fig. 2c).

137 Another way to ascertain the extent of past changes is through the inspection of historical time se-
138 ries of direct temperature measurements or documentation of such environmental observations. The 20th
139 century showed a series of decadal-scale climate anomalies such as the 'Great Salinity Anomaly' (GSA)
140 observed in the subpolar North Atlantic in the late 1960s and the early 1970s (Dickson et al., 1988) or
141 the Early Arctic Warming (EAW) in the 1920s and 1930s (Schokalsky, 1936; Bengtsson et al., 2004). The
142 EAW might be seen as a regime shift in the North Atlantic ocean and atmosphere responsible for a re-
143 gional warming (Bengtsson et al. 2004; Drinkwater, 2006; Tokinaga et al., 2017), whereas the GSA can
144 be linked to a sea ice transport out of the Arctic through Fram Strait (Hilmer et al., 1998; Häkkinen &
145 Geiger, 2000), which represents a major source of freshwater for the northern North Atlantic (Aagaard
146 & Carmack, 1989; Schmith & Hansen, 2003). The GSA was accompanied by a subsequent decadal cool-
147 ing in the 1970s over the North Atlantic realm (Dima & Lohmann, 2007). These events could be consid-
148 ered as analogues to past Heinrich events which appeared during glacial times (Heinrich, 1988), whereas
149 the EAW has a flavor of a DO-type warming over Greenland. The knowledge about potential drivers and
150 the inherent dynamics is still poorly known. We show a potential link between synoptic atmospheric vari-
151 ability of days and decadal-to-multidecadal time scale variability, e.g., the GSA, droughts and floods, ex-
152 treme winters, and atmospheric blocking (Dickson et al., 1988). Increased sea ice export from the Arc-
153 tic Ocean stabilizes the upper water column in the North Atlantic, diminishes the production of inter-
154 mediate and deep water masses through ocean convection and can influence the large-scale ocean circu-
155 lation (Häkkinen, 1999). An important part of sea ice export from the Arctic is forced by specific atmo-
156 spheric structures (Hilmer et al., 1998; Häkkinen & Geiger, 2000; Cavalieri 2002; Ionita et al., 2016). These
157 can be parts of coupled ocean-atmosphere modes of climate variability with large-scale projections.

158 Examples of abrupt transitions and extremes are known for the last century, and the multi-decadal
159 variability is coupled to the persistence of atmospheric blocking and synoptic variability (Pfahl et al., 2009;
160 Rimbu & Lohmann, 2010, 2011; Rimbu et al., 2014, 2016a). Paleoclimate data allows to evaluate climate
161 variability and extremes and their relation to the mean climate. A model-based approach will help to get
162 a more consistent picture to obtain a common model-data interpretation of the underlying mechanisms
163 and forcings of paleoweather events in the North Atlantic/European realm. High-resolution environmen-
164 tal archives bring the relatively short period of the instrumental record into a long-term context. Tran-
165 sitions can be detected with the highest possible resolution which can range from decadal up to seasonal
166 or even synoptic time scales. We further propose future research directions for studying climate changes
167 regarding sampling strategies, statistical tools, and the formulation of climate models, reviewing the chal-
168 lenges to understand the causes and contributors to climate shifts and extremes.

169 **2 Temperature signature due to orbital variations**

170 A fundamental feature of Earth’s climate during the last few million years is the regular fluctua-
171 tion between different states. The driving effect is due to the geometry of the Earth’s orbit affecting the
172 latitudinal and seasonal distribution of insolation (Milankovitch, 1941; Berger, 1979; Laskar et al., 2004).
173 By computational analysis of the planetary system, these variations can be calculated to a high accuracy
174 for the last millions of years (Laskar et al., 2004). Apart from shaping the seasonal cycle, the changes in
175 the seasonal and latitudinal distribution of insolation are a primary driver for climate variability on multi-
176 millennial time scales (Huybers & Curry, 2006). This relationship has been hypothesized for a long time
177 (Adhemar, 1842; Croll, 1875), and coherent variability of climate records and orbital parameters were found
178 when geologists started to date climatic proxy records of ocean sediments (e.g., Broecker & van Donk,
179 1970; Hays et al., 1976). One classical concept for the response of the climate system on orbital forcing
180 was proposed by Imbrie et al. (1992) in that insolation changes at high northern latitudes initiate the
181 climate response. However, the question remains how the climate system reacts to the local insolation
182 forcing.

183 To quantify the response of Earth’s climate to orbital forcing during the last 1 My (My= 10^6 years),
184 we produce a new numerical simulation of transient climate that is forced with the solution of orbital el-
185 ements by Laskar et al. (2004). We employ the Community Earth System Models (COSMOS) framework

186 to provide a numerical representation of the coupled atmosphere-ocean-land system, which is described
187 in Appendix A. In order to compute monthly mean climate over the last 1 My, we prescribe a time se-
188 ries of eccentricity, obliquity, and precession of the Earth’s orbit. These values are sampled at a time-resolution
189 of 100 years, which results in the acceleration of orbital forcing by a factor of 100. Such an acceleration
190 method is necessary to derive the numerical solution of 1 My of climate in a coupled atmosphere-ocean
191 general circulation model within a reasonable amount of time.

192 Our analysis is based on simulated changes in near surface air temperature (SAT), that reflects a
193 direct response to orbitally modulated spatial and seasonal distribution of insolation at the top of the
194 atmosphere. We study the climate’s response to forcing on orbital time-scales. Hence, our focus is in par-
195 ticular on those modes of climate variability that are linked to the impact of precession, obliquity, and
196 eccentricity of the Earth’s orbit around the Sun. We perform a spectral analysis of simulated transient
197 near-surface temperature over the last 1 My and compute the integrated spectral power density in three
198 different periodicity bands: precession band (18 ky - 26 ky), obliquity band (38 ky - 56 ky), and eccen-
199 tricity band (90 ky to 110 ky). The 420-ky orbital periodicity is not analyzed here because we concen-
200 trate on the last 1 My. A fast Fourier transform is applied on detrended time series at grid resolution of
201 the atmosphere model across the entire globe. We provide three different graphical representations of the
202 power spectral density for the last 1 My. We illustrate the latitudinal dependency of the zonally aver-
203 aged power spectrum (Fig. 4). This illustration highlights the substantial dependence of the strength of
204 the climate’s response to latitude and periodicity of orbital forcing. The frequencies of precessional re-
205 sponse are detected at ~ 23 and ~ 19 ky, together with a first harmonic with half of the years lined to the
206 semi-precession cycle with much lower amplitude. The precessional response is strongest in the subtrop-
207 ics, but also visible at polar latitudes in the Northern Hemisphere and at high latitudes in the Southern
208 Hemisphere. The precessional response is related to nonlinearities and/or seasonal biases in the climate
209 system. Furthermore, we see power in the obliquity band (41 ka, 29 ka and 54 ka), mostly at high lat-
210 itudes with a small maximum in the subtropics.

211 In order to see the frequency bands explicitly, geographical maps highlight regions that are partic-
212 ularly sensitive to a specific periodicity of orbital forcing (Fig. 5). We provide the power spectrum den-
213 sity over periodicity and latitude, which is done for the entire surface of the planet as well as for land and
214 ocean separately. Fig. 5a shows that the precessional response dominates in subtropical Africa, India,

215 South America, and some regions around Greenland. Obliquity (Fig. 5b) shows a similar pattern, but
216 with more emphasis on the region around Greenland and suppressed response across South America as
217 compared to precession. The high latitude response is not that different over land and ocean. The 100
218 ky-band (Fig. 5c) emphasizes again the subtropical region in northern Africa and some regions in the trop-
219 ics and mid-latitudes, but with a much smaller amplitude than precession and obliquity. The changes in
220 eccentricity of the Earth's orbit have a small effect on the global mean insolation but modulate the strength
221 of the seasonality and therefore have an effect through the precession. The 100-ky frequency band in part
222 simply originates from an amplitude cycle of precession cycles, a linkage that may control the similar-
223 ities shown by the local power spectrum of annual mean SAT (Fig. 5a and c).

224 The frequency analysis is also applied to other climate phenomena like the ocean circulation (not
225 shown). The wind-driven ocean circulation is dominated by obliquity (more a linear response), whereas
226 the Atlantic meridional overturning and heat transport show pronounced precessional and obliquity peaks
227 due to nonlinearities and seasonal biases towards winter. Our model integrations and analysis confirm
228 that convection and deepwater formation serve to rectify the zero annual-mean precessional forcing, re-
229 sulting in precessional energy in the ocean. The Antarctic Bottom Water also shows eccentricity peaks
230 indicating rectification mechanisms.

231 The local approach complements the global Milankovitch/Imbrie hypothesis (climate is driven by
232 northern summer insolation) in explaining observed climate variability and potentially offers new insights
233 in interpreting paleoclimate records, especially beyond the 100 ky responses. It shows the limitations of
234 global calibration curves based on Northern Hemisphere insolation changes at high latitudes. The pro-
235 nounced high-latitude response both in precession and obliquity emphasizes the important feedbacks and
236 non-linearities including the ocean circulation (deep water formation is happening in winter) and land
237 (ice-albedo feedback). Actually, the ice-albedo feedback is underrepresented in our long-term simulation,
238 as ice sheets are static. The climate-vegetation feedback has so far not been analyzed this in the 1 Ma
239 simulation. Interesting is the variability in the subtropics. The insolation forcing favors a north-south
240 seasonal march of the ITCZ (Intertropical Convergence Zone) and the inland penetration of Afro-Asian
241 monsoon precipitation during boreal summer. Surface temperature is reduced in regions where precip-
242 itation is enhanced due to the combination of increased cloud cover and increased surface evaporation
243 (Herold & Lohmann, 2009; Braconnot et al., 2019). Braconnot et al. (2019) demonstrate that the pro-

244 jection of temperature onto the insolation curve is larger in the Northern than in the Southern Hemisphere,
245 especially between 10-40°N where about 80% of the temperature signal is a direct response to the local
246 insolation forcing. A completely different approach is followed by the template model (Laepfle & Lohmann,
247 2009) where insolation-driven temperature variability on orbital time scales relies on the modern rela-
248 tionship between insolation and temperature throughout the year. Over extratropical continental areas,
249 a linear response is detected, whereas over the ocean, the subtropics, and sea-ice sensitive regions, pre-
250 cessional bands at 19 and 23 ky are visible related to nonlinearities such as a seasonal mixing regime. Such
251 an empirical approach includes all fast feedback processes from the seasonal time scale. In a similar way,
252 the climate sensitivity can be constrained from the seasonal cycle in temperature (Knutti et al., 2005).
253 Qualitatively, COSMOS and the template model show similar patterns. Of course, the full climate model
254 has "longer-term" feedbacks like the ice/snow-albedo feedback on millennial time scales and non-local
255 effects giving higher responses at most latitudes. However, there are also clear differences. The full cli-
256 mate model shows a strong response in the obliquity band in northern high latitudes, which is absent in
257 the template model without long-term feedbacks. Our results are different to the results from Short et
258 al. (1991), who used a linear two-dimensional energy balance model to study the spatially resolved tem-
259 perature response of the last 800 ky. However, as their model was linear, precession and eccentricity only
260 appeared when using the maximum temperature as diagnostics. Our study shows that the climate archives
261 recording annual mean temperature will also reflect these frequencies caused by non-linearities in the cli-
262 mate model already.

263 **3 Millennial time-scale variability**

264 Dansgaard-Oeschger and Heinrich events are pronounced climatic changes over the last 120,000 years.
265 Although many of their properties were derived from climate reconstructions, the associated physical mech-
266 anisms are not yet fully understood. We are interested in the exact position in time where a change in
267 temperature occurred on the given time scale. According to different DO-criteria we determined several
268 sequences of DO-onsets and are interested in any potential periodic recurrence. A standard technique would
269 be a spectral analysis of the detrended time series and the identification of significant peaks in the fre-
270 quency spectrum (e.g., Grootes & Stuiver, 1997). Here, we focus on the pacing of discrete DO-onsets which
271 is the reason why we apply a different measure of periodicity based on Huybers und Wunsch (2005) and

272 Ditlevsen et al. (2007). It allows us to evaluate the phase stability of the discrete event onsets although
273 we have a small sample size (Huybers & Wunsch, 2005). This of course also applies for quasi-periodic sys-
274 tems with an unpredictable component where events repeat themselves irregularly or are skipped. In Ap-
275 pendix B, we list the timing of the DO-events for the last glacial (Table 2). There, we describe the thresh-
276 old approach where we consider change points when a certain threshold criterion is fulfilled. To quan-
277 tify such uncertainties, we furthermore provide a probabilistic methodology and use a Bayesian approach
278 to identify the DO-events in time. The algorithm in use for this study was provided by Ruggieri (2013)
279 and has the advantage of quantifying uncertainties in number and location of warming events (Appendix
280 B). Fig. 6 shows the $\delta^{18}O$ signal of the NGRIP ice core from Greenland covering the last glacial-interglacial
281 period covering the past 120 ky and change-point statistics.

282 Repeatedly, a periodicity of about 1500 years is reported in the literature, starting with the spec-
283 tral analysis of the GISP2 and GRIP by Grootes and Stuiver (1997) with a peak at 1470 years (Yiou et
284 al., 1997). Also with a focus on the recurrence interval of partially skipped warming events in the Green-
285 land record an inter-event waiting time of 1470 years is announced (Alley et al., 2001; Schulz, 2002; Rahm-
286 storf, 2003). Ditlevsen et al. (2007) include also the Bølling/Allerød period and event 9 which was re-
287 ported by Rahmstorf (2003) after applying a low pass filter and defining a 2‰ anomaly in a 200 yrs time
288 interval as a DO-criterion. Rahmstorf (2003) supported the idea of regular paced discrete events by a pe-
289 riod of 1470 years, by showing that 13 of his defined events fall at least in an interval of 20% of multi-
290 ples of this period. Here we present a simple stochastic process given random input for the occurrence
291 of DO events (Ditlevsen et al., 2007). Consequently, we are testing the inter-event waiting times in obey-
292 ing the exponential distribution which is by definition of a Poisson Process (Kroese et al., 2011) and can
293 formally be stated as our null hypothesis.

294 We will show that the analysis is statistically robust with regard to the exact methodology for event
295 identification as we base our analysis on the dating reported in previous studies (Ditlevsen et al., 2007;
296 Barker et al., 2011; Rasmussen et al., 2014) (cf. Table 2). Using Monte Carlo simulations we generate
297 1000 randomly sampled time sequences of events with a waiting time distribution. We then calculate the
298 test statistics of each Monte Carlo simulation which estimates the distribution under our null hypoth-
299 esis and then tests whether the score of the data (different DO-time sequences) lies within a 90% con-
300 fidence interval, corresponding to a conservative testing. The Rayleigh R measure is chosen because it

301 captures information about the regularity which is in question here (Mardia & Jupp, 2009). For unper-
 302 turbed periodic event occurrence, it results in a maximum value of 1. In contrast to Ditlevsen et al. (2005),
 303 we will also report the periods for which we obtain the 5 most dominant peaks (Table 1). We evaluate
 304 the test for the whole period (Table 1a) and the 42-10 ky BP interval (Table 1b) of the NGRIP record.
 305 For the shorter interval, all R values for major periods are greater than 0.5. The prominent 1470 yrs pe-
 306 riod results in a $R_{max} = R(P1)$ for the 2 thresholded sequences and is rated as the second-highest pe-
 307 riod for the other DO-sets. For Barker et al. (2011) and Rasmussen et al. (2014) data sets, the domi-
 308 nant peak in the Rayleigh R test is for 940 years. Generally, Barker et al. (2011) (bark11) and Bayes
 309 (bayes) result in higher R values because fewer events are considered. The Bayes sequence attains its R_{max}
 310 for a period of 548 years, with respect to the high inter-event waiting time of 3908 years ("Estimate" in
 311 Table 1b) which would mean that the recurrence is highly irregular and several periodic beatings are skipped.
 312 Regardless of the period analysed, the measure for periodicity for all sequences falls within a high like-
 313 lihood region of a stochastic driven process such that we cannot reject our null hypothesis of statistical
 314 random occurrence of DO-events. What is interesting here is that the periods 1470, 900, and 1150 yrs
 315 (colored red, blue, green in Table 1) also obtain very high Rayleigh R values, which fall well in a 90% con-
 316 fidence interval for the probability density from exponentially distributed waiting times. This is because
 317 the Rayleigh R values for the dominant periods for NGRIP (42-10 ky) do not deviate significantly from
 318 the true R for random event time sequences with exponentially distributed waiting times.

319 One additional circumstance that needs to be contemplated is the small sample size. Ditlevsen et
 320 al. (2007) only considered a maximum number of 12 events. During the most recent glacial cycle we con-
 321 sider only 18-29 events which might be one reason why we find those striking periodicities (high R val-
 322 ues). If DO-events were really drawn from a random process then we would expect a decreasing Rayleigh
 323 R (decreasing measure for periodicity) with increasing sample size. That is why we are also interested
 324 in periodic components on a longer time scale. Although the abrupt warmings are best dated and researched
 325 on a shorter interval (50-10 ky) they are encountered throughout the whole last glacial and beyond. We
 326 see that the Rayleigh R values decrease which is consistent with our null-hypothesis. Interestingly, the
 327 periods ~ 1500 , ~ 1150 years do not dominate the range of periods which result in high R values, more
 328 outstanding is the period of ~ 900 years (Table 1a). To overcome the problem of small sample size, one
 329 could also perform the analysis for the predicted DO-event occurrence of the synthetic Greenland record

of Barker et al. (2011). The difficulty is then that the dating becomes more uncertain. Low R values with $R_{max} < 0.4$ indicate a weak periodicity for all periods when using Barker et al. (2011). High-resolution and accurate dating seems to be necessary to get a more coherent view of the full system. Similar statistical data analysis is necessary to build hypotheses related to events on decadal and centennial time scales. In the following, we will examine some of the major challenges in Earth System Models to be applied for long-term climate and weather changes.

4 Methodological challenges of Earth System Models

The continuous evolution of climate models over recent decades has been enabled by a considerable increase in computational capacity, with supercomputer speeds increasing by more than a factor of 10^6 since the 1970s (IPCC, 2007). This computational progress has permitted a corresponding increase in model complexity (by including more and more Earth system components and relevant processes), in the achievable length of the simulations, and in spatial resolution (Claussen et al., 2002). To significantly improve the simulation of climate by general circulation models, systematic errors in representations of relevant processes must first be identified, and then reduced. This endeavor demands that the parameterizations of unresolved processes, in particular, should be tested over a wide range of time scales (weather, climate, **and** paleoclimate).

4.1 Internal variability: challenges in models

Different flavors of DO-type oscillation are reported for models of different complexity (Winton, 1993; Yin & Sarachik, 1995; Wang & Mysak, 2006; Kim et al., 2012; Peltier & Vettoretti, 2014; Li & Born, 2019). The underlying causes responsible for millennial-scale variability are still debated. Ice core data suggests a strong coupling of the largest millennial-scale warm events in Antarctica with the longest DO-events in Greenland through the AMOC (Barbante et al., 2006; Dima et al., 2018). A quasi-global impact of the events has been reported in the literature (Voelker, 2002). Marine data suggest an anti-phase correlation between AMOC and Atlantic subsurface temperatures under glacial climate conditions, suggesting that the vertical temperature structure and associated changes in AMOC are a key element governing DO cycles (Kim et al., 2012; Sadatzki et al., 2019). As intermediate ocean layers are gaining buoyancy as a result of downward mixing of the warmer overlying water, a point is reached at which water

357 at the surface becomes denser than those in the mid-depth and deep ventilation is initiated (Kim et al.,
358 2012). Sea ice expansion preceded the buildup of a deep oceanic heat reservoir (Sadatzki et al., 2019).
359 During a transition to a strong AMOC, former sites of maximum convection are still there, and the Labrador
360 Sea and the North Atlantic open ocean appear as additional sites of deep convection. During the strong
361 mode, the North Atlantic basin gets colder by convection and fresher until the heat import from lower
362 latitudes can not overcome the freshening anymore. Thus, the strong convection is inhibited and the weak
363 mode of the oscillations is reactivated. In a recent comprehensive review (Li and Born, 2019), the authors
364 examine evidence that DO events are an unforced or spontaneous oscillation of the coupled atmosphere-
365 ice-ocean system comprising the subpolar gyre as a key region producing DO-like events. Interestingly,
366 similar behavior is reported under pre-industrial conditions in a long control integration (Sidorenko et
367 al., 2015) (Fig. 7a). In a pre-industrial control setup these events could be judged as undesirable. Yet,
368 ultimately this result shows value in suggesting that on the one hand current models include the physics
369 that are necessary to produce abrupt climate transitions, but on the other hand - as Li and Born (2019)
370 emphasized - exhibit incorrect sensitivity to the boundary conditions. Through our time series analysis,
371 we have now another constraint to the models, namely that the oscillations should be irregular with ran-
372 dom waiting times. On shorter time scales Danek et al. (2019) find that the resolution in the ocean cir-
373 culation model matters for the Labrador Sea convection (Fig. 7b). The high-resolution run shows gen-
374 eral agreement to observations (Good et al. 2013). A basic feature is a sporadic deepening of the mixed
375 layer depth, related to turbulent processes on the meso- and sub-mesoscale within the Labrador Sea in-
376 terior. As one driver, we identify the atmospheric circulation with strong winds over the Labrador Sea
377 area (Fig. 7c). More research is required to study the physical processes affecting the mixing at these sen-
378 sitive sites, including sub-mesoscale turbulence (Callies et al., 2015). It seems that irregularities on mil-
379 lennial as well decadal time scale are an important feature of the convection and thus of the large-scale
380 overturning circulation, again bridging time and spatial scales (Fig. 1).

381 In further model developments (Sidorenko et al., 2018), the large oscillations in the Labrador Sea
382 mixing were reduced compared to the old experiments. However, it might be that the centennial-to-millennial
383 oscillations are required to explain climate variability as expressed e.g. by the Little Ice Age and the Me-
384 dieval Warm Period during the last 1000 years (Bradley, 2014; Broecker, 2000). Recent model develop-
385 ment is mainly done with a particular focus on the present and future climate. It could be that a de-tuning

386 of the models is necessary in order to exhibit irregular oscillations on centennial-to-millennial time-scales.
387 A systematical analysis of the mechanisms leading to centennial-to-millennial variability remains open.
388 Numerical experiments (Danek et al., 2019) suggest that at least in the Labrador Sea and other sensi-
389 tive areas high resolution can play an important role in realistically simulating the variability in the mixed
390 layer depth affecting AMOC.

391 Besides the vertical dipole, there seems a bi-polar signature of the DO events like for the termina-
392 tion (Broecker, 1998; Steig & Alley, 2002). During the first stage of the DOs, Antarctica steadily warmed,
393 but little change occurred in Greenland. Thereafter, at the time when Greenland’s climate underwent
394 an abrupt warming, the warming in Antarctica stopped and reversed. During the final phase, a sudden
395 decrease in the northward heat transport cools the north (Dima et al., 2018). CO₂ variations are an ac-
396 tive part in the system on millennial time scales. Changes are dominated by slow changes in the deep ocean
397 inventory of carbon, correlated with Antarctic temperature and Southern Ocean stratification, while AMOC
398 changes are more responsible for fluctuations on millennial time scales (Schmittner & Galbraith, 2008).
399 Moreover, carbon dissolved in the global ocean may also be coupled to geosphere-climate interactions (Hasen-
400 clever et al., 2017; Wallmann et al., 2018).

401 4.2 Deglacial meltwater and climate

402 Sea-level reconstructions indicate a massive input of meltwater during the early stage of the last
403 deglaciation. During the Bølling period (14.7–14.1 ky before present) up to ~ 0.5 Sv ($1 \text{ Sv} = 10^6 \text{ m}^3$
404 s^{-1}) of meltwater entered the ocean (Fairbanks et al., 1992). Most of the meltwater likely originated from
405 the decaying Laurentide ice sheet; the inferred freshwater pulse out of the Mississippi River entered the
406 Atlantic via the Gulf of Mexico (Fairbanks et al., 1992). Various modeling studies indicate that such mas-
407 sive meltwater input should lead to the cessation of NADW, a reduction in large-scale AMOC, and a cool-
408 ing (e.g., Stocker et al., 1992; Rahmstorf, 1996; Manabe & Stouffer, 1997; Kageyama et al., 2013). A sig-
409 nificant drop in AMOC is, however, not supported by paleoclimate data (Sarnthein et al., 1994; McManus
410 et al., 2004). In a quite different approach for the last deglaciation, Knorr & Lohmann (2007) illustrated
411 that the onset of a strong AMOC is inevitable, and that only the timing of AMOC resumption is trig-
412 gered by melting water, indicating that the freshwater flux plays only a secondary role for the termina-
413 tion of ice sheets. On the contrary, a more recent study by Liu et al. (2009) emphasizes a huge influence

414 of freshwater perturbation to the AMOC history where a negative freshwater input was applied to get
415 a warm Bølling period.

416 Sensitivity experiments to high-latitude freshwater perturbations started with climate model sce-
417 narios of the polar halocline catastrophe by Bryan (1986). As stability and response of the AMOC to iden-
418 tical freshwater perturbation scenarios vary greatly between different models (e.g., Manabe & Stouffer,
419 1997; Rahmstorf, 1996; Kageyama et al., 2013), we concentrate here on the dynamics of the deglacial melt-
420 water in climate. Again, by applying our multi-resolution approach in the Finite Element Sea ice Ocean
421 Model FESOM (Wang et al., 2014, Sidorenko et al., 2015), we will be able to zoom into regions of inter-
422 est (coasts, high latitudes, equator) while keeping the resolution sufficiently low in other areas (Fig. 8).
423 The model includes furthermore a land surface scheme with interactive vegetation dynamics, ensuring
424 that climate and vegetation are consistent with each other. The model is described in the Appendix C.

425 Before exploring the effects of different regions of freshwater forcing, re-call that deglacial freshwa-
426 ter fluxes are by iceberg calving and surface runoff from the ice sheets. Such meltwater fluxes were sim-
427 ulated by the ice sheet model of Zweck & Huybrechts (2005), shown in Fig. 9. For the period 14.7–14.1
428 ky BP, most meltwater entered as liquid water, whereas during the Heinrich-1 event (ca. 17.5–14.8 ky
429 BP) the freshwater entered as solid water in terms of calving. The calved icebergs bring the freshwater
430 very effectively into the subpolar region (Heinrich, 1988), and it is a common approach to mimic this fresh-
431 water input as a broad anomaly into the subpolar North Atlantic Ocean (e.g., Rahmstorf et al., 2005).

432 We start now with water hosing experiments based on a glacial background state which has been
433 obtained after 840 years of model integration. After the spinup, we insert a 150-year meltwater flux of
434 0.2 Sv perturbation at different locations (Mississippi, Mackenzie, or local coastal runoff close to Labrador).
435 Significant differences are found in the oceanic responses to these different locations of freshwater discharge
436 (Fig. 10a). The AMOC-index refers to the maximum of the AMOC in the North Atlantic Ocean. The
437 Mississippi (orange) and the Mackenzie (blue) runoff have little influence, in contrast to the subpolar forc-
438 ing (red) where the freshwater flux is equally distributed in the 50 – 70°N-band in the North Atlantic
439 Ocean. As an additional experiment, the hosing is inserted at the coast close to the Labrador Sea between
440 50 – 70°N (green). The different response in the AMOC can be attributed to the negative sea surface
441 salinity anomalies lowering the density (Fig. 10b). The low salinities prevent deep water formation and
442 can lead to a partial collapse of the AMOC, especially in the subpolar setup (red).

443 The weakened overturning circulation further suppresses heat exchange between the North Atlantic
444 surface and sub-surface water masses, leading to a cooling in the upper ocean and a pronounced warm-
445 ing in the sub-surface layer (Fig. 11). The signature bears similarities with the anti-phase correlations
446 between AMOC and Atlantic subsurface temperatures for the DO events. Again, the intermediate ocean
447 is de-densified by the downward mixing. The resumption of the deep ventilation (Fig. 10a) is the stronger,
448 the more effective the subsurface warming is (Fig. 11). In the case of the subpolar hosing, the overshoot
449 is around the model year 1100. Note, that the smallest response and the least efficient hosing is in the
450 Gulf of Mexico, a potential region for a major MWP-1a discharge (Aharon, 2003, 2006). The results are
451 in line with previous uncoupled ocean model experiments (Condron & Winsor, 2011) showing that the
452 narrow coastal boundary currents play an important role efficiently preventing the freshwater entering
453 into the deep-water formation areas.

454 Fig. 12 shows the near-surface air temperature anomaly for the different glacial and deglacial sce-
455 narios. The temperature response at the Mackenzie, Mississippi, and coastal scenarios are relative insen-
456 sitive compared to the subpolar hosing. Interestingly, the subpolar shows a pronounced overshoot after
457 the hosing, releasing heat stored at the subsurface (Fig. 11). Fig. 13 shows the associated streamfunc-
458 tion of AMOC, indicating the strong changes for the subpolar scenario both during the hosing and in the
459 overshoot phases. The AMOC in other scenarios is rather insensitive to the freshwater forcing. As an al-
460 ternative explanation for the insensitivity of the Bølling freshening, it has been suggested that the ocean
461 circulation is too sensitive to surface freshening when underestimating the overflow contribution of the
462 deep water seeping from the Nordic Seas (Lohmann, 1998), and that a most likely northern source of deep
463 water formation can stabilize the system during the Bølling (Lohmann & Schulz, 2000). Another pos-
464 sibility follows the approach of Roche et al. (2007), where the deglacial meltwater pulse could have sneaked
465 unnoticed into the deep ocean during the last glacial due to the effect of hyperpycnal flow. The exact dy-
466 namics of the AMOC perturbation is not elaborated here and is subject of a subsequent analysis. It can
467 be shown that connections between latitudes are due to boundary-trapped Kelvin waves propagating along
468 western boundaries. For interannual or larger periods, the mid-latitude Kelvin waves are replaced by long
469 Rossby waves on the eastern boundary and viscous boundary waves on the western boundary. The lo-
470 calization of the viscous boundary waves is on the so-called Munk boundary layer scale which is much
471 larger than the Rossby radius (Kawase, 1987; Johnson & Marshall, 2002, 2004).

472 Source and causal mechanism of MWP-1a are still being explored. Based on corals drilled offshore
473 from Tahiti, Deschamps et al. (2012) show that MWP-1a started no earlier than 14,650 years ago and
474 ended before 14,310 years ago, making it coeval with the Bølling, but most likely including a significant
475 meltwater contribution from the Southern Hemisphere (see also Weber et al., 2014, for a discussion). Gre-
476 goire et al. (2012) found in their ice-sheet modeling simulation that the separation of the Laurentide and
477 Cordilleran ice sheets in North America produces a meltwater pulse corresponding to MWP-1a. Stan-
478 ford et al. (2006) derive that MWP-1a did not coincide with the sharp Bølling warming, but instead with
479 the abrupt cooling of the Older Dryas with only a relatively minor 200-year weakening of NADW flow,
480 which is very much in line with our findings.

481 Summarizing, we argue that the way how meltwater enters into the ocean (liquid or solid) provides
482 a first order control for the sensitivity of the ocean circulation and thus the deglacial climate. For future
483 studies, a fully coupled climate–ice sheet model capturing the drainage chronology of the Laurentide Ice
484 Sheet is being developed to shed more light on the role of deglacial meltwater on climate. A high reso-
485 lution at the coasts seems to be an important feature for a proper understanding of the coastally-confined
486 freshwater pathways. Indeed, using a coarse-resolution version of the Condrum & Windsor (2011)-model,
487 it was demonstrated that the spreading of freshwater across the subpolar North Atlantic results from the
488 inability of numerical models of this resolution to accurately resolve narrow coastal flows, producing in-
489 stead a diffuse circulation that advects freshwater away from the continental boundaries. In order to cap-
490 ture the climatic impact of fresh-water released in the past and possible future, the ocean needs to be mod-
491 eled at a resolution sufficient to resolve the dynamics of narrow, coastal buoyant flows. Another difficulty
492 in resolving the history of abrupt climate change during the last glacial period is beyond the sensitivity
493 of climate models to freshwater and to the location of the freshwater source. It is directly related to the
494 dynamic behaviour of the ice sheets. To this end, proper modeling of the ice-substrate interface is an im-
495 portant boundary condition for ice sheet modeling. The substrate affects the ice sheet by allowing slid-
496 ing through sediment deformation and accommodating the storage and drainage of subglacial water (e.g.,
497 Gowan et al., 2019). Temperate ice sheets, such as the Laurentide and Eurasian ice sheets behaved dif-
498 ferently depending on whether or not there were thick, continuous unconsolidated sediments underneath
499 the ice (Clark & Walder, 1994), providing an uncertainty in the ice sheet history of deglacial climate.

4.3 Radiocarbon dynamics for chronology

Marine climate records are frequently dated by means of stratigraphic approaches such as orbital tuning or the synchronization with other cross-dated geological records. A more physical approach is radiocarbon (^{14}C) dating which takes advantage of the half-life of radioactive ^{14}C (5700 years, Audi et al. 2003; Bé et al. 2013). Radiocarbon is produced through cosmogenic radiation in the upper atmosphere from where it enters the global carbon cycle as $^{14}\text{CO}_2$. Cosmogenic ^{14}C production has been varying (e.g., Hain et al. 2014; Adolphi et al. 2018; Channel et al. 2018), and in addition, the partitioning of ^{14}C between the various carbon reservoirs has been subject to past climate-carbon cycle interactions (e.g. Köhler et al. 2006; Hesse et al., 2011).

The ^{14}C uptake by the oceans is slow (~ 10 years) and superimposed by mixing with water from the ^{14}C -depleted deep sea (Broecker & Peng, 1974). Therefore, marine ^{14}C concentrations are systematically depleted with respect to ^{14}C in the overlying contemporaneous atmosphere. The concentration difference translates into an apparent surface water ^{14}C age (or "marine reservoir age") ranging from about 300 years in the subtropics to up to 1000 years in high latitudes during the late Holocene (e.g. Key et al., 2004). Marine reservoir ages can have even a much broader range of 100 - 2500 years at sites far offshore (e.g., Skinner et al., 2010; Cook & Keigwin, 2015; Sarnthein et al., 2015).

A further complication is that marine ^{14}C and ^{14}C -dated climate records typically originate from continental margins, marginal seas, or tropical islands (see Fig. 14a). Such sites may not be representative to mirror large- or global-scale processes of climatic change, demanding a valid extrapolation method to infer past global conditions. The spatial gaps can be filled employing ocean-climate circulation models. However, the current generation of ^{14}C -equipped models suffers from insufficiently low resolution near the ocean margins (e.g., Butzin et al. 2017), and conventional high-resolution ocean modeling approaches are computationally too expensive when it comes to long-term simulations covering the entire time scale accessible to ^{14}C dating (\sim the past 55 ky). This problem can be solved by means of global multiresolution models using unstructured meshes (Appendix C). Fig. 14 shows a first simulation result for the late Holocene obtained with a ^{14}C -equipped version of the global multi-resolution model FESOM2 (Danilov et al., 2017). With this approach we will be able to zoom into regions of interest while keeping the resolution sufficiently low in other areas (Fig. 8). By means of conventional ocean models, it is difficult to resolve important interbasin exchange fluxes through narrow key seaways like the Bering Strait, Strait

529 of Gibraltar, Indonesian Throughflow, and Canadian Arctic Archipelago. Unstructured meshes have the
530 potential to represent such throughflows with higher accuracy.

531 A major challenge in paleoclimatology is to disentangle the relationship between climatic leads and
532 lags during the past 50 ky (e.g., Shakun et al., 2012). In the case of radiocarbon dating this implies to
533 quantify the uncertainties associated with climatic variability and cosmogenic ^{14}C production changes,
534 and that the time axis might be changed (Fig. 3). Here, we consider an ensemble of simulated marine
535 surface water ^{14}C concentrations (Fig. 15), using a more efficient model version (Butzin et al., 2005, 2017).
536 The atmospheric ^{14}C forcing is based upon the Hulu Cave speleothem record including uncertainties (Southon
537 et al. 2012; Cheng et al. 2018; Butzin et al., 2020). We combine three different ocean climates and ^{14}C
538 forcing scenarios to create an ensemble of nine experiments. Each experiment is spun up over 20 ky with
539 fixed atmospheric values of ^{14}C and CO_2 at 54 ky before present, and is then run transiently forced by
540 time-variable ^{14}C (Fig. 15a) and CO_2 . Marine reservoir ages for the upper 50 m are calculated afterwards
541 (Appendix C). Radiocarbon concentrations in terms of the fractionation-corrected $^{14}\text{C}/^{12}\text{C}$ ratio of ma-
542 rine surface water normalized to the corresponding ratio in the atmosphere ($F^{14}\text{C}$) are shown in Fig. 15b.
543 Fig. 15c shows the ^{14}C age averaged between 50°N - 50°S , and Fig. 15d a snapshot for the Last Glacial
544 Maximum at 21 ky BP. Ensemble approaches provide a new direction in evaluating dating uncertainties
545 on regional scale and a dynamical interpretation of paleoclimate data (Alves et al., 2018), and thus sub-
546 stantially enhance our understanding of the Earth's climate system.

547 **5 Variability of extreme climate and weather over Europe**

548 In the atmosphere, preferred modes of climate variability at different time scales exist, which of-
549 ten span large geographical areas. Some regions may be cooler or perhaps drier than average, while at
550 the same time, thousands of kilometers away, warmer and wetter conditions prevail. These simultane-
551 ous variations in climate, often of opposite sign, over distant parts of the globe, are commonly referred
552 to as teleconnections in the literature (e.g., Barnston & Livezey, 1987; Trenberth et al., 1998; Wallace &
553 Gutzler, 1981). Examples are the Atlantic Multidecadal Oscillation (AMO), Pacific Decadal Oscillation
554 (PDO), El Niño-Southern Oscillation (ENSO) (cf. Fig. 1). The teleconnections are an important tool
555 for the interpretation of proxy data, as they link a time series at one point to large-scale patterns. Be-
556 low, we will discuss the mechanisms behind decadal to multidecadal variability and extreme climate dur-

557 ing the recent past and their possible dynamics (e.g., Pfahl et al., 2009; Rimbu & Lohmann, 2010, 2011;
558 Rimbu et al., 2014, 2016a; Otto, 2017; Sippel et al., 2020). Our analysis is based on advanced statisti-
559 cal techniques applied to observational data, high-resolution proxy records, and climate simulations. Ex-
560 treme climate events are typically associated with specific weather patterns and coherent atmospheric
561 structures, e.g. cyclones and blockings, but also to large-scale and long-term variations. We show again
562 a linkage of the dynamics across time scales.

563 **5.1 Decadal time scales: the Great Salinity Anomaly**

564 Beyond freshwater variations during the last deglaciation, that had an impact on circulation pat-
565 terns in the North Atlantic Ocean as outlined above, changes in North Atlantic freshwater have also been
566 observed for recent climate. Major freshening episodes have been described in recent decades (Curry &
567 Mauritzen, 2005). The first freshening event has happened in the 1970s, the so-called the Great Salin-
568 ity Anomaly (GSA) (Dickson et al., 1988). Interestingly, the GSA which was accompanied with a decadal
569 cooling over the North Atlantic realm, might have some analogues to past Heinrich events which appeared
570 during glacial times (Heinrich, 1988). The GSA in the subpolar North Atlantic in the 1970s, is also linked
571 to droughts and floods, extreme winters, and atmospheric blocking. Aagaard and Carmack (1989) reported
572 that the origin of this event must have been an anomalously large Arctic freshwater discharge through
573 Fram Strait, later confirmed by Häkkinen (1993) using a coupled ocean-sea ice model. The 1970s GSA
574 was followed by similar events (but weaker in amplitude) in the 1980s (Belkin et al., 1998) and the 1990s
575 (Belkin, 2004). The events have also been linked to atmospheric variability (Haak et al., 2003; Karcher
576 et al., 2005; Dima & Lohmann, 2007; Mauritzen et al., 2012).

577 Fig. 16 indicates a potential link of synoptic variability to decadal-to-multidecadal time scale vari-
578 ability (Ionita et al., 2016). During times of enhanced blocking between Greenland and Svalbard (e.g.
579 in 1962-1966), sea ice first accumulated in the Arctic, later leading to an enhanced sea ice release into
580 the North Atlantic, and with a delay reducing Labrador Sea salinity. In the following decade, the fresh-
581 ening causes a weakening of the large-scale AMOC. The AMOC goes to a weaker state several years af-
582 ter the GSA in the late 1960s, consistent with earlier finding of a fast response of the meridional over-
583 turning to North Atlantic forcing (Eden & Willebrand, 2001). From this perspective, the 1960s GSA in-
584 duces an AMOC shift to a new weaker state caused by the significant fresh water forcing related to Fram

585 Strait Sea Ice Export (Fig. 16). For more recent times, a close link between decadal atmospheric block-
586 ing is proposed (Häkkinen et al., 2011), but knowledge about the inherent atmosphere-ocean-sea ice dy-
587 namics and their potential drivers is still incomplete. When comparing the three major freshwater events
588 in the Northern North Atlantic (Fig. 16c), the Fram Strait Sea Ice Export and atmospheric blocking are
589 not the only causes for large-scale ocean circulation changes on these time scales.

590 Decadal-to-multidecadal variability is coupled both to synoptic time scales, such as atmospheric block-
591 ing and long-term background conditions. As an example, the jet stream in the Euro-Atlantic region and
592 associated blocking activity is linked to Atlantic multidecadal ocean variability and AMOC, implying the
593 need for a proper representation of synoptic scale variability in coupled climate models. Model simula-
594 tions are advancing towards increased resolution, the addition of feedback mechanisms between differ-
595 ent components of the earth/climate system, and, in particular, the added ability of considering the dy-
596 namic simulation of isotopes. Future research will use a hierarchy of global climate simulations and all
597 available types of observational, (paleo-)environmental and reanalysis data to evaluate variability pat-
598 terns and extreme events on synoptic to millennial time scales. This integrated approach provides a long-
599 term context beyond our limited view from the observational period. The combined analysis of data and
600 models is essential for interpreting paleoclimate data and related weather and climate patterns, which
601 is the topic of the next subsection.

602 **5.2 Changes in the intensity and frequency of extremes: Atmospheric rivers**

603 Changes in the intensity and frequency of extremes have drawn much attention during recent decades
604 (IPCC report on Managing the Risks of Extreme Events and Disasters to Advance Climate Change Adap-
605 tation, 2012; Coumou et al., 2014; Horton et al., 2016), mainly due to their large impacts on natural en-
606 vironment, economy and human health (Ciais et al., 2005; Kovats & Kristie, 2006). Due to the inherent
607 nature of extreme events, the variability of high-temperature extremes differs from that of mean temper-
608 ature (Raible et al., 2007; Schär, 2015). Therefore, we need to understand the causes of climate extremes
609 in the 20th century and in a more distant past. Paleoclimate data and modelling allow an assessment of
610 climate variability and extremes and their relation to the average climate.

611 One particular example of extreme events is linked to atmospheric rivers (AR) which are narrow
612 regions responsible for the majority of the poleward water vapor transport across the midlatitudes. They

613 are characterized by high water vapor content and strong low level winds of extratropical cyclones (e.g.,
614 Namias, 1939). Newell et al. (1992) termed these long (~ 2000 km) and narrow (300–500 km wide) bands
615 of enhanced water vapor flux tropospheric rivers (Fig. 17). They used the term rivers because they trans-
616 port water at volumetric flow rates similar to those of the world’s largest rivers, through narrow corri-
617 dors (Zhu & Newell, 1998). Although the meridional water vapor transport within ARs is critical for wa-
618 ter resources, ARs can also cause disastrous floods especially when encountering mountainous terrain (Pasquier
619 et al., 2018). Here, we describe future direction and research needs of the synoptic interpretation of pa-
620 leoclimate data.

621 Rimbu et al. (2016b) show that lake sediment records from Ammersee (southern Germany) are a
622 good proxy not only for local summer extreme precipitation. The significant peaks of flood frequency (Fig.
623 18a) characterize also extreme local precipitation variability as well as extreme high temperatures over
624 northeastern Europe. River Ammer floods are associated with enhanced moisture transport from the At-
625 lantic Ocean and the Mediterranean towards the Ammer region, a pronounced trough over western Eu-
626 rope as well as enhanced potential vorticity at upper levels. Interannual to multidecadal increases in flood
627 frequency are associated with a wave train pattern extending from the North Atlantic to western Asia,
628 with a prominent negative center over western Europe. This link between low and mid-latitudes is ex-
629 plained by the atmospheric circulation pattern associated with AR (Fig. 17b). It has been shown that
630 a large proportion of the most intense precipitation events (and of course of their associated floods) in
631 western Europe are objectively associated with the occurrence of ARs, both in Great Britain (Lavers et
632 al., 2012) and in the Iberian Peninsula (Ramos et al., 2015). As ice cores, corals, as well as lake sediments
633 go back in time from centuries to several millennia with high, i.e. annual or even seasonal resolution, we
634 can use them to reconstruct past variability of related climate extremes. This is useful not only for knowl-
635 edge of past extreme climate variability, but also to anticipate future natural variability of climate ex-
636 tremes.

637 **6 Conclusions and Outlook**

638 The Earth’s climate is characterized by many modes of variability in the atmosphere, ocean, cryosphere
639 and biosphere. The generation of low-frequency variability in the climate system is crucial to allow a sep-
640 aration of anthropogenic signals from natural variability, thereby increasing the ability to recognize and

641 improve the attribution of climate and weather events to anthropogenic climate change. In the paper,
642 we first consider the response of the system to regular multi-millennia orbital forcing, and then consider
643 shorter time scales down to daily weather events (Fig. 1).

644 Milankovitch (1941) initially suggested that the critical factor for ice sheet formation and decay is
645 total summer insolation at about 65°N , because for an ice sheet to grow some additional ice must sur-
646 vive each successive summer. In contrast, the Southern Hemisphere is limited in its response because the
647 expansion of ice sheets is curtailed by the Southern Ocean around Antarctica. Thus, the conventional view
648 of glaciation is that low summer insolation in the temperate Northern Hemisphere leads to perennial con-
649 tinental snow triggering Northern Hemisphere glaciation. Here, we have evaluated how the climate sys-
650 tem creates non-linear responses to **local insolation forcing**. Our model results show a pronounced po-
651 lar amplification of the orbital forcing. In the frequency domain, temperature (and precipitation) vari-
652 ability at mid-latitudes are dominated by precession, while high latitudes are dominated by obliquity. The
653 precessional response is related to nonlinearities and/or seasonal biases in the climate system. Additional
654 to the insolation forcing, greenhouse gases provide a net forcing driving glacial-interglacial variability. The
655 order of magnitude of glacial-interglacial changes is similar in magnitude to our current anthropogenic
656 induced radiative forcing of $\sim 2 \text{ W m}^{-2}$, which will increase in the future. Surface temperatures are cur-
657 rently not in equilibrium due to the thermal inertia of the ocean, and the Earth's present energy imbal-
658 ance of $0.5\text{--}1 \text{ W m}^{-2}$ (von Schuckmann et al., 2016) provides an indication of how much additional global
659 warming is still in the pipeline if climate forcings remain unchanged (Hansen et al., 2016). We can learn
660 about potential thresholds in the system from the past (e.g., Kwasniok & Lohmann, 2009; Robinson et
661 al., 2012; Sutter et al., 2016). As an example, the Mid-Pleistocene Transition at about 1 My marks a tran-
662 sition from a 40-ky variability to a dominant 100-ky variability in the climate system. Lowering of global
663 CO_2 below some critical threshold, conceivable triggers a transition from a 40 ky- to a 100 kyr-resonance
664 (e.g., Raymo et al., 1998; Tziperman et al., 2006; Farmer et al., 2019). Further insight may come from
665 the analysis of the oldest ice in Antarctica, providing information of CO_2 and Southern Hemisphere tem-
666 peratures (Fischer et al., 2013), as well as through modeling of ice sheets that are truly coupled into Earth
667 system models.

668 We have questioned regularities found in DO-events occurrence and statistically analysed the dis-
669 tribution of inter-event waiting times. Our assumption was that DO-events represent repeated oscilla-

tions between interstadial and stadial conditions with a **threshold crossing process** involved (Alley et al., 2003; Ditlevsen et al., 2007; Steffensen et al., 2008). That assumption is not only based on the general shape of DO-events as rapid transitions between binary states, but also concluded from model simulations (Ganopolski & Rahmstorf 2001) that examined the stability of the glacial climate with the usage of ocean-atmosphere model of intermediate complexity. This has been recently confirmed in a more comprehensive model (Zhang et al., 2017). It was reported that DO-like events could be caused by a threshold-driven process, as a result of changes in the Atlantic freshwater balance combined with the ice sheet height, consistent with the data analysis of Schulz et al. (1999). Therefore, we have considered a threshold criterion for the definition of DO events and used a Bayesian methodology for accurate change-point detection that allows us to quantify the uncertainties about the total number of change-points and their locations (Ruggieri, 2013; Partin et al., 2015) . For a given time interval each event occurs stochastically independent, meaning that the probability of one abrupt warming event does not affect the probability of any other warming event in the same interval. We find that the observed DO-recurrence intervals in the NGRIP ice core proxy data are consistent with exponentially distributed waiting times not only in the interval of 42 -11 kyrs BP (as proposed by Ditlevsen et al., 2007), but also during the whole last glacial period in the last 120 ky. This finding is independent from the dating methodology. Additionally, novel periodicities of ~ 900 and ~ 1150 years in the NGRIP record are reported besides the prominent 1500-years cycle. Although a high periodicity reflected in a high Rayleigh R can be found in the data, it remains indistinguishable from a simple stationary random Poisson process. These are quite remarkable findings as ~ 1500 and ~ 900 yrs periods are visible throughout the Holocene (Bond, 1997; Rimbu et al., 2004; Xu et al., 2015).

Another feature for millennial climate variability comes from the spatio-temporal structure of DO events and modeling results. DO-like variability is seen in a handful of model simulations, including even some pre-industrial simulations. As a mechanism, the subsurface is warmed by the downward mixing of the warmer overlying water during an AMOC weak state, until the surface becomes more dense than at mid-depth and deep ventilation is initiated. A similar mechanism is observed in the overshoot dynamics of AMOC.

Decaying Northern Hemisphere ice sheets during the last deglaciation affected the high latitude hydrological balance in the North Atlantic and hence the ocean circulation after the Last Glacial

699 Maximum. Surprisingly, geological data suggest that meltwater fluxes of about 14-20 m sea-level equiv-
700 alent flushed into the North Atlantic without significantly influencing the Atlantic meridional overturn-
701 ing circulation. Global sea level reconstructions indicate abrupt changes within several hundred years and
702 rising the sea level during the Bølling/Allerød warm interval (Clark et al., 1996, 2002a). Using a climate
703 model with high resolution near the coast, we have investigated the response of the ocean circulation to
704 deglacial freshwater discharges. In our experiments, we find a strong sensitivity of the ocean circulation
705 depending where the deglacial meltwater is injected. Meltwater injections via the Mississippi and near-
706 Labrador coast hardly affect the AMOC. The reduced sensitivity of the overturning circulation against
707 freshwater perturbations following the Mississippi route provides a consistent representation of the deglacial
708 climate evolution. However, a subpolar freshening of the North Atlantic, mimicking a transport of wa-
709 ter by icebergs, leads to a quasi-shutdown of AMOC. Therefore, during Heinrich events the freshwater
710 injection is more effective. Future model developments shall therefore concentrate on modeling of ice calv-
711 ing (Alley et al., 2008; Levermann et al., 2012) and iceberg transport in the ocean (Rackow et al., 2018b;
712 Stern et al., 2019). Most ocean climate models do not represent ice shelf calving in a physically realis-
713 tic way (and even not properly parameterized), although calving of icebergs is a major component of the
714 mass balance for present Antarctic ice shelves. It is remarkable that the vertical heat distribution mat-
715 ters for the AMOC dynamics. The strongly-weakened AMOC (as in the case of subpolar hosing) leads
716 to upper ocean cooling and pronounced warming in the sub-surface layer (Rühlemann et al., 2004). In
717 this particular scenario, we see an AMOC overshoot, also seen during the Bølling/Allerød. At millennial
718 scales, paleoceanographic evidences have characterized abrupt changes in the other regions as well, e.g.
719 the subpolar North Pacific Ocean (Rae et al., 2014; Maier et al., 2018; Lembke-Jene et al., 2018; Lohmann
720 et al., 2019; Gong et al., 2019). Such rapid shifts affect processes like surface stratification, formation and
721 dynamics of sea ice, feedbacks with atmospheric patterns like the Aleutian Low, as well as deep and in-
722 termediate water formation across time scales. Studies of abrupt events are of great interest because they
723 could help to indicate whether rapid reorganizations of the ocean circulation might occur in the future,
724 and how they might affect climate in general (Thirumalai, 2018).

725 One of the greatest obstacles to resolve the leads and lags-puzzle is the difficulty of developing an
726 accurate time scale for long ice cores and marine sediments. The development of such a time scale would
727 allow testing of many climate-forcing hypotheses, leading to significant advances in paleoclimate theory.

728 In the late Quaternary, **the backbone of such a time scale could be provided through radiocar-**
729 **bon dating.** To improve the quality of the interpretation of the available radiocarbon records, high spa-
730 tial model resolution is required around the relevant locations, which makes traditional ocean-climate mod-
731 els with their uniform meshes impractical. Here, traditional high-resolution modeling approaches, involv-
732 ing uniform meshes, would result in prohibitive computational costs. To overcome this drawback, one can
733 apply a coupled climate-radiocarbon circulation model using innovative and scalable numerical schemes
734 on supercomputing facilities (e.g., Koldunov et al., 2019). This issue deserves further efforts combining
735 simulations and reconstructions, all the more, as uncertainties in reconstructions are large for times prior
736 to about 30 ky before present, where the marine reservoir variability is poorly constrained through re-
737 constructions. Data assimilation approaches could provide a physically consistent climate-carbon isotope
738 history. The data assimilation methodology combines the data (observations, reconstructions) and the
739 underlying dynamical principles governing the system to provide a state estimate of the system which
740 is better than what could be obtained using just the data or the model alone. Applying an ensemble Kalman
741 filter approach, one could use the advances in a parallel data assimilation framework with relatively mod-
742 erate increase in computation time (Nerger & Hiller, 2013).

743 As changes at high-latitudes are particularly sensitive to forcing, the detection and understanding
744 of polar-to-subpolar changes and their influence on global climate is of central importance. Model sim-
745 ulations can aid the interpretation of the drivers and causes of observed variations in paleoclimate quan-
746 tities. **Analyzing proxy-reconstructed paleoclimate records and models in tandem allows for**
747 **the evaluation of forcing and feedback mechanisms under climate change.** Conversely, model
748 simulations can aid in the interpretation of the causes of observed variations in paleoclimate data. One
749 important application of paleoclimate data is to validate state-of-the-art coupled climate models for past
750 climate transitions, which allows to gain insight into possible future climate states that may be notably
751 different from present day conditions. Past time periods provide the means for evaluating the performance
752 of general circulation models for interglacials (e.g. Lohmann et al., 2013a; Shi & Lohmann, 2016; Pfeif-
753 fer & Lohmann, 2016; Stein et al., 2017). Key findings are that the models do not capture the magni-
754 tude of the sea surface temperature anomalies derived from marine proxy records for interglacials, and
755 that they underestimate the variance (Laepple and Huybers, 2014a,b). The connection between climate
756 sensitivity and variance is consistent with a possible connection of the linear response to perturbations

757 and the fluctuations (Lohmann, 2018). Indeed, Nijssen et al. (2019) emphasized in a similar way that mod-
758 els that are more sensitive to greenhouse gas emissions (i.e. higher equilibrium climate sensitivity) also
759 have higher temperature variability on time scales of several years to several decades (Lenton, 2011). There-
760 fore, high-sensitivity climates are more prone to rapid warming but also to hiatus periods.

761 Observed and simulated statistics of extreme events shall be compared between different model con-
762 figurations and observations to evaluate if, and under what conditions climate models are reliable for sim-
763 ulating future changes in climate extremes. **Through a combination of paleoclimate records with**
764 **the analysis of climate-model simulations, we can contribute to a better understanding of**
765 **the long-term evolution of climate variability modes and their linkage to synoptic events.**
766 Such strategy can overcome our current knowledge gaps, with special emphasis on the role of ocean cir-
767 culation, climate modes of variability, and regional extreme climatic events. Numerical climate models
768 operate on increasingly finer grid sizes as the performance of parallelized super computers increases. Whether
769 a model can represent a geophysical process depends on the model formulation and discretization. Since
770 the spatial scale of oceanic eddies is $O(1-100)$ km (first baroclinic Rossby radius of deformation), the ocean
771 model grid resolution needs to be on the same scale to represent these ubiquitous small-scale features (Chel-
772 ton et al., 2011). Alternatively, their effects must be parameterized as in current state-of-the-art general
773 circulation models, where the typical horizontal resolutions of the ocean component is about 1 degree.
774 Practically, given the present high-performance computer capacities, efficient and parallelized model codes,
775 it is now possible to conduct isotope-enabled (e.g., Werner et al., 2018; Cauquoin et al., 2019) simula-
776 tions for 50-100 model years per day even with a multi-scale ansatz (Shi et al., 2020).

777 **Recent developments have considerably improved the computational efficiency and scal-**
778 **ability of unstructured-mesh approaches on high-performance computing systems** (Danilov
779 et al., 2017). The surface ocean current in such high-resolution simulation (Fig. 19) has a completely dif-
780 ferent structure including eddies than the structure in coarse-resolution model. Conceptual and statis-
781 tical approaches are used to explore climate variability from the persistence of atmospheric blocking to
782 non-linear responses to external forcing to multi-millennial variability. Sea ice decline has a potential strong
783 influence on the mid-latitude climate. Results from recent modeling experiments indicate that the influ-
784 ence of Arctic sea ice decline is not sufficiently strong to explain anomalously cold winters in Europe, but
785 that other regions such as Russia and North America may be more strongly affected than Europe (Semm-

786 ler et al., 2016). The uncertainty surrounding modeled Arctic-mid-latitude linkages has recently led to
787 the Polar Amplification Model Intercomparison Project (Smith et al., 2019) which could be extended to
788 past intervals which would be a step beyond current activities in the paleoclimate modeling community
789 (Kageyama et al., 2017; Otto-Bliesner et al., 2017).

790 The simulated interglacial periods provide the means for evaluating the performance of general cir-
791 culation models in representing sea surface temperature (SST) anomalies and trends (e.g., Lohmann et
792 al., 2013a; Pfeiffer & Lohmann, 2016; Stein et al., 2017). However, the models do not capture the mag-
793 nitude of the derived SSTs from marine proxy records in all climate simulations of the Holocene (Lohmann
794 et al., 2013a) and Last Interglacial (Pfeiffer & Lohmann, 2016). It is suggested that a part of such dis-
795 crepancies can be caused either by too simplistic interpretation of the proxy data (including dating un-
796 certainties and seasonal biases) and/or by underestimated regional responses in climate models. By us-
797 ing long-term multi-millennial climate model runs and paleoclimate data, a discrepancy with a too low
798 low-frequency variance is detected also with respect to variability (Laepfle & Huybers, 2014a,b). The
799 model-data differences on long time scales suggest that relevant feedback mechanisms and internal vari-
800 ability may not be well represented in current climate models (Lohmann, 2018).

801 One of the critical feedback mechanisms for long-term variability is the interaction of Arctic sea ice
802 and the global ocean circulation, which can affect climate variability on short time scales and might lead
803 to relatively rapid climate change (Häkkinen et al., 2011; Ionita et al., 2016). Improving our understand-
804 ing of the sea-ice feedback, Müller and Stein (2014) reconstructed the sea ice conditions in the Fram Strait
805 at the end of the last glaciation at centennial to millennial time scales. For more recent times, sea ice can
806 accumulate in the Arctic due to atmospheric blocking over Greenland (Ionita et al., 2016), leading finally
807 to an enhanced sea-ice release into the North Atlantic reducing Labrador Sea salinity affecting AMOC.
808 It is likely that new model approaches are necessary to realistically obtain persistent blocking in conjunc-
809 tion with large-scale changes in the ocean circulation and sea ice (Kinter III et al., 2013). **We suggest**
810 **a close link between climate variability modes and the statistics of weather patterns in the**
811 **North Atlantic realm. Such features provide a challenge of numerical models and the in-**
812 **terpretation of long-term climate records.** We provide an example of atmospheric rivers and ex-
813 treme events in Europe as evaluated from paleoclimate data (Rimbu et al., 2016b). Atmospheric rivers
814 account for nearly all of the extreme flooding events with potentially huge economic losses (e.g., Dominguez

815 et al., 2018). The potential effects of extreme events, such as heat waves or droughts, do not only depend
816 on their abundance, but also on "how these extremes occur", including the sequence of events. These events
817 are unexplored for both past and future changes, and they require sophisticated analysis methods (Sedlmeier
818 et al., 2016). Weather and climate extremes are an inherent part of climate and variability on synoptic
819 time scales can have a large effect on decadal-to-multidecadal variability and vice versa.

820 Looking at the development of science in the 20th century one can distinguish two trends, which
821 Weisskopf (1963) called "intensive" and "extensive" research. Intensive research goes for the fundamen-
822 tal laws, extensive research goes for the explanation of phenomena in terms of known fundamental laws.
823 There is always much less intensive research going on than extensive. Once new fundamental laws are
824 discovered, a large and ever increasing activity begins in order to apply the discoveries to hitherto un-
825 explained phenomena. **Climate science is mostly extensive and most fundamental laws were**
826 **discovered decades ago** (Landau & Lifshitz, 1987)(although there are Navier-Stokes existence and smooth-
827 ness problems, see: The Clay Mathematics Institute, <http://www.claymath.org/millennium-problems/>).
828 However, many modeling and methodological issues of obtaining new data are related to current inten-
829 sive research or at least intensive research of yesterday. Examples are recent breakthroughs in compu-
830 tational fluid mechanics (e.g., Danilov et al, 2017) and in high-precision radiocarbon dating (e.g., Wacker
831 et al., 2010). Kuhn (1962) argues in terms of the evolution of scientific theories that different phases do
832 not emerge from the straightforward accumulation of facts, but rather from a set of changing intellectual
833 circumstances and possibilities. Climate science already went through a pre-paradigm phase in the sense
834 of Kuhn (1962), in which there is no consensus on any particular approach. Somehow, climate and gen-
835 eral Earth System Science are partly in a second phase already, in which the open questions are solved
836 within the context of the established methods documented in review articles, books and reports (e.g., Berger,
837 1988; Peixoto & Oort, 1992; Marshall & Plumb, 2007; Holton & Hakim, 2012; McGuffie & Henderson-
838 Sellers, 2014). However, note that it has been argued that ocean carbon models might be still in the il-
839 lusion and pre-paradigm phase (Le Quere, 2011). From our discussions above we conclude that more sci-
840 entific efforts yield to more complex views, which can be even contradictory. This indicates that at least
841 some underlying assumptions might need to be revisited, reflecting phase three of Kuhn's (1962) theory
842 of scientific development. **In modeling, there is a clear "dilemma of growth" in that the mod-**
843 **els are required for understanding the dynamics of a system as complex as the Earth Sys-**

844 **tem.** On the other hand, the growing complexity of the models themselves seems to jeopardize under-
845 standing (Held, 2005; Gramelsberger et al., 2020). Therefore, different levels of complexity are required
846 to understand the Earth system on long time scales (e.g., Timmermann & Lohmann, 2000; Ganopolski
847 & Rahmstorf 2001; Claussen et al., 2002; Tziperman et al., 2006; Lohmann & Ditlevsen, 2019).

848 Future research may be enhanced along three directions data, statistics, and models, leading to an
849 increase in the current knowledge about the climate evolution. From an epistemological viewpoint, it is
850 crucial that researchers deepen or acquire the ability to integrate all three directions into their arsenal
851 of methods. Since it is hard to disagree with the simple statement "more data are better", the task here
852 is rather to identify those dimensions in the data space where invested resources may yield to a maxi-
853 mum of new information. In this way, data assimilation techniques could help for an estimate of the state
854 of the system, but also its uncertainty (Kalman, 1960; Burgers et al., 1998; Nerger & Hiller, 2013). Fi-
855 nally, a detailed analysis of paleoclimatology shall test the assumptions about the climate recorder sys-
856 tems regarding leads, lags, and other filter properties (Laepfle et al., 2011; Lohmann et al., 2013b).

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857
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863 for this research are included in this paper are NGRIP (2004) with the time scale of Andersen et al. (2007),
864 Köhler et al. (2017), Archer & Brovkin (2008), Stepanek and Lohmann (2020), Butzin et al. (2012b), Sidorenko
865 et al. (2015), Danek et al. (2019), Zweck & Huybrechts (2005), the ¹⁴CHRONO Marine Reservoir Database
866 (calib.org/marine), using Reimer et al. (2004, 2013) and further references therein. Further datasets used
867 in this paper are NCEP reanalysis data (Kalnay et al., 1996), Czymzik et al. (2010). Figure 15 uses data
868 from Southon et al. (2012), Cheng et al. (2018), Skinner et al. (2017) and Butzin et al. (2019). Figure
869 16 is based on Ionita et al. (2016). Figure 18 is based on Czymzik et al. (2010) and Rimbu et al. (2016b).
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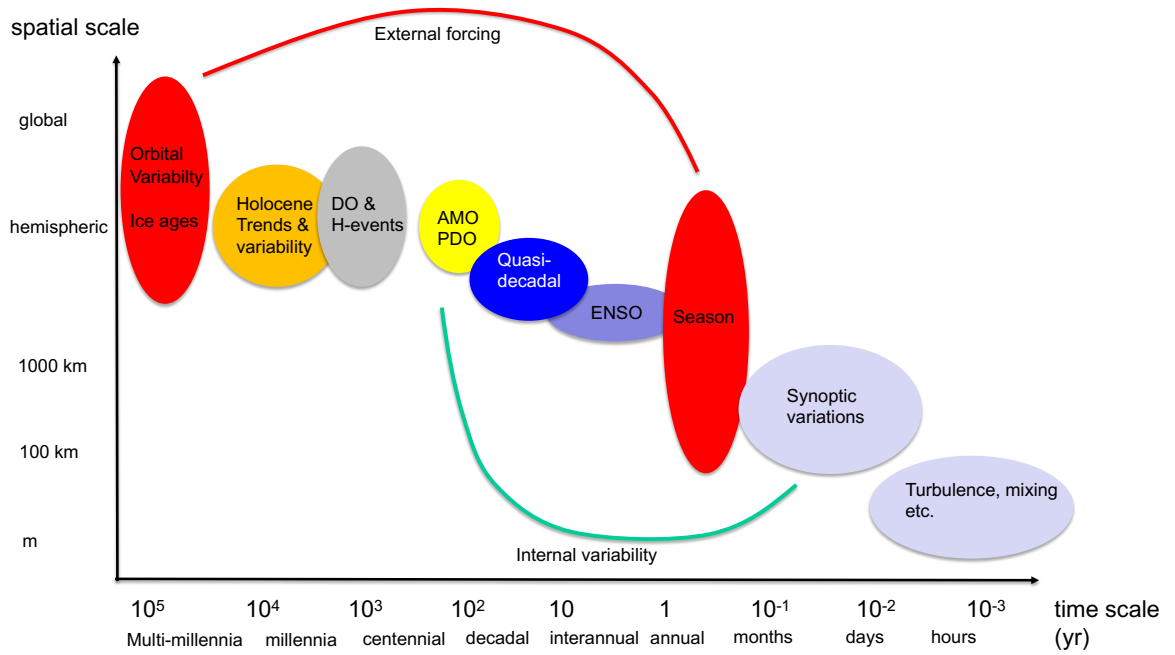


Figure 1. Schematic diagram of the spatio-temporal scales considered. DO: Dansgaard-Oeschger, H: Heinrich events; AMO: Atlantic Multidecadal Oscillation; PDO: Pacific Decadal Oscillation; ENSO: El Niño-Southern Oscillation. The annual and astronomical cycles are externally driven and have quasi-global impact.

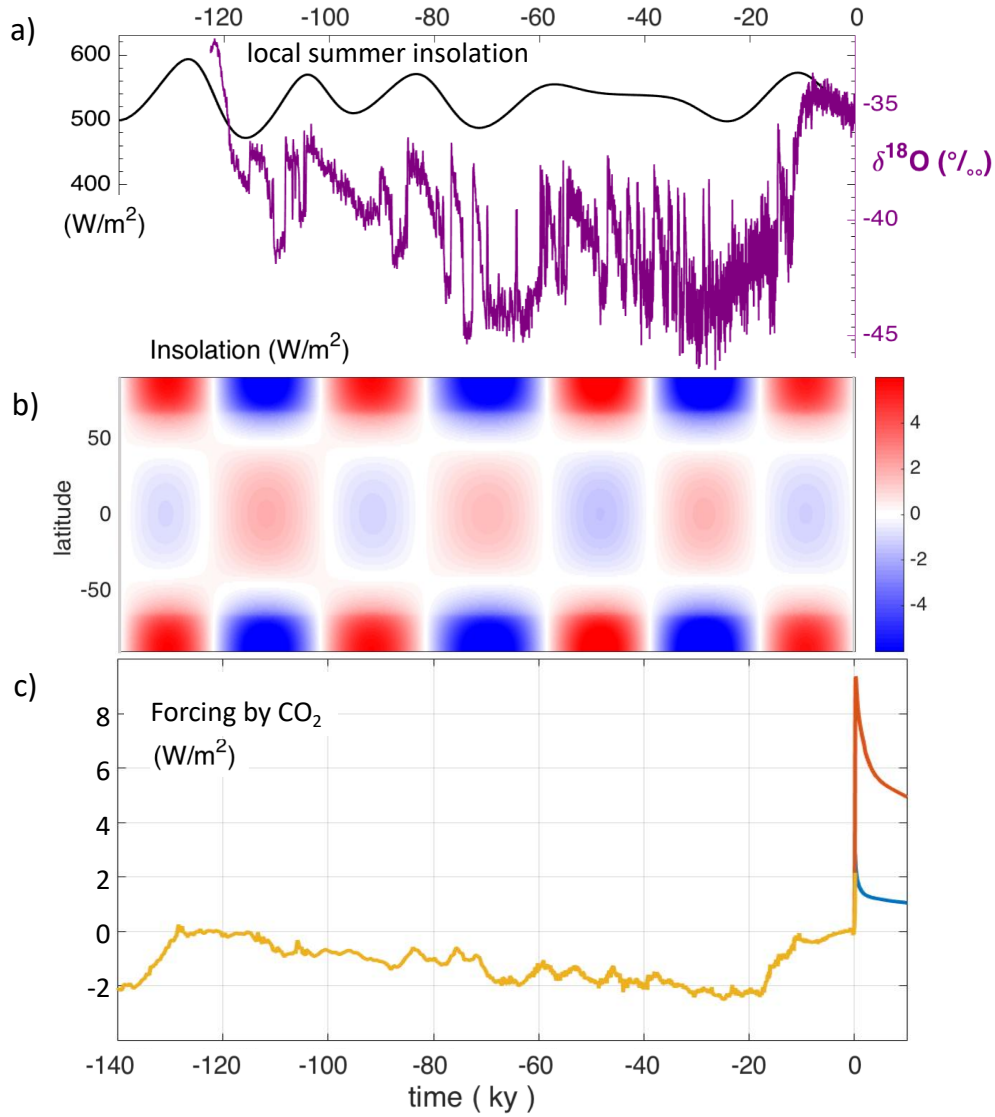


Figure 2. a) The $\delta^{18}\text{O}$ ice core curve from the North Greenland Ice Core Project (NGRIP, 2004) documents climate variability over the last 120 ky (purple). The black curve indicates the June 21 insolation at the local position (Wm^{-2}). b) Annual mean insolation variation at all latitudes using the algorithm of Berger (1978). c) CO_2 -forcing as reconstructed from the past (yellow) (Köhler et al., 2017) and estimated for future scenarios (Archer & Brovkin, 2008) for a moderate (1,000 Gt carbon) (blue) and large (5,000 Gt carbon) (red) fossil fuel slugs (the natural atmospheric CO_2 content is on the order of 600 Gt carbon prior to anthropogenic combustion of carbon). For the translation into Wm^{-2} , we assume a 4 Wm^{-2} for doubling of CO_2 and a logarithmic dependence $\ln(\text{CO}_2/\text{CO}_2^{\text{ref}})$ with CO_2 and CO_2^{ref} being the CO_2 -level and the reference pre-industrial level, respectively.

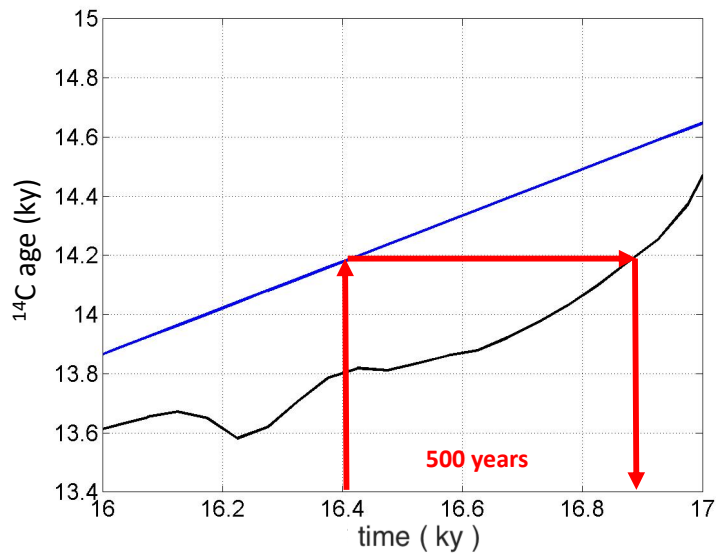


Figure 3. Reservoir age correction procedure to convert from ^{14}C years to calendar years. The blue line uses the Reimer et al. (2004) calibration, the black line uses output of a climate-radiocarbon model (Butzin et al., 2012a,b). The mismatch at a hypothetical data point at 16.4 ky is about 500 years.

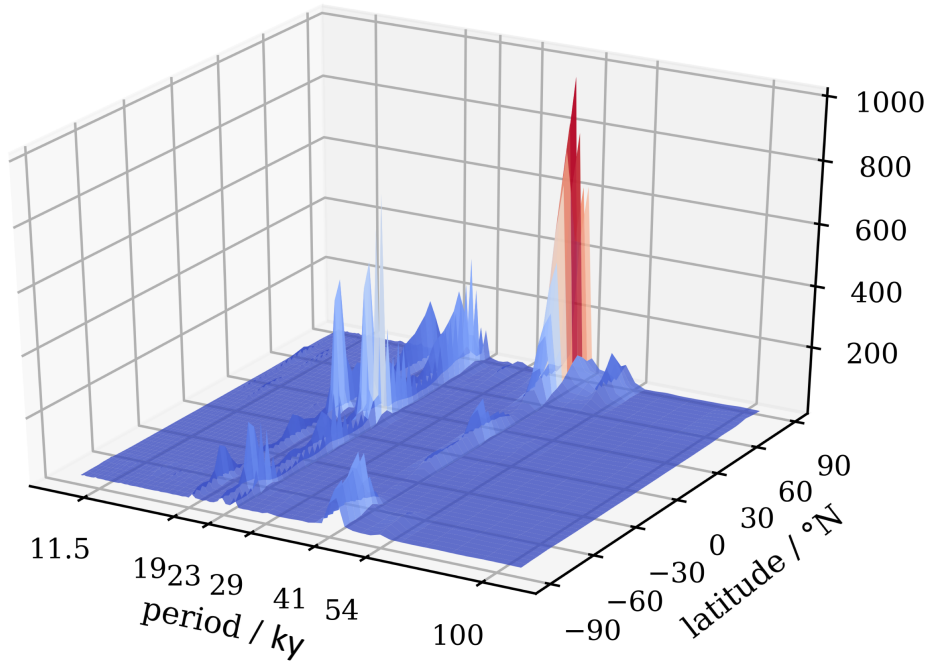


Figure 4. Power spectrum density of zonal mean annual mean surface air temperature (SAT) as calculated from the climate model COSMOS.

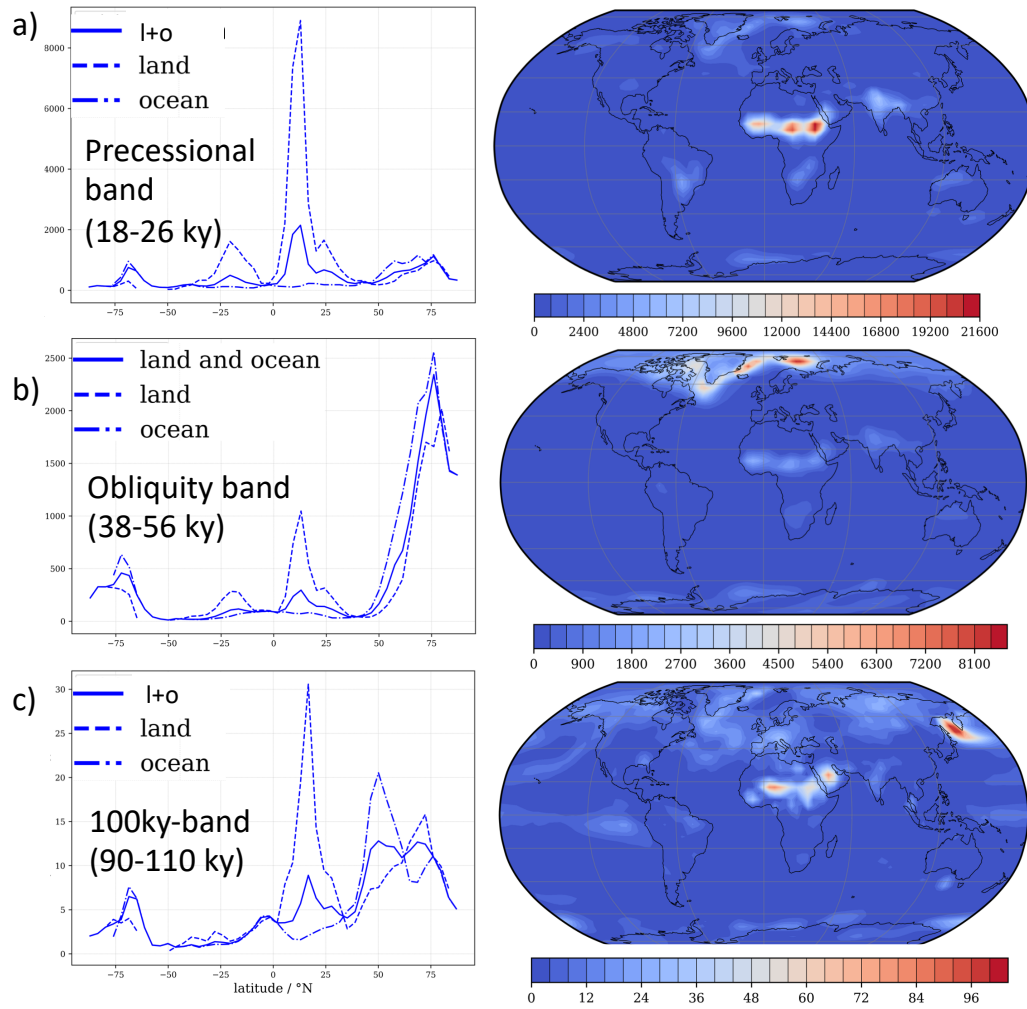


Figure 5. Power Spectrum Density for different frequency bands of annual mean SAT. Left: Zonal mean with solid lines for land and ocean, dashed line for land, dashed-dotted lines for the ocean. Right: local power spectrum of annual mean SAT. Note the different scales. The precessional band is dominated by a peak over North Africa (a), obliquity for a region around Greenland and the Barents Sea (b), eccentricity with a much lower amplitude over North Africa and the Sea of Okhotsk (c).

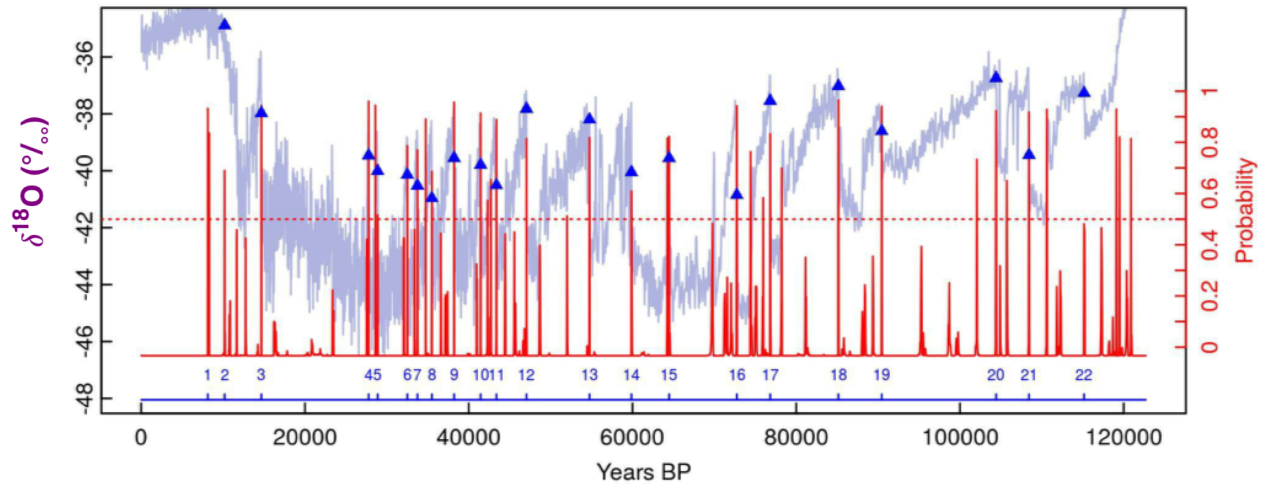


Figure 6. Bayesian Change Point Analysis for the $\delta^{18}O$ signal of the NGRIP ice core from Greenland based on GICC05 time scale (Andersen et al., 2007). Note that time evolves from right to left in accordance to geological convention. Red spikes indicate the probabilities for a change point at each time point proposed by the Bayesian Change Point Algorithm (Ruggieri, 2013). Warming and cooling change points are detected. Blue triangles mark the warming events with a probability higher then 0.5 (dotted red line), numbered at the bottom of the graph and used for further analysis.

a) Past 120 ky

Seq.	# Ev.	Estimate	P1	P2	P3	P4	P5	
ras14	26	4147	649	596	1282	1799	519	yr
			0.487	0.463	0.427	0.413	0.397	R
bark11	22	4526	871	573	1305	967	763	yr
			0.622	0.477	0.453	0.452	0.446	R
thresh	24	4369	827	956	1020	546	510	yr
			0.563	0.522	0.482	0.474	0.47	R
bayes	21	5248	517	4446	958	637	1380	yr
			0.525	0.499	0.459	0.443	0.438	R

b) Period of 10-42 ky

Seq.	# Ev.	Estimate	P1	P2	P3	P4	P5	
ras14	11	2976	948	1488	594	1574	649	yr
			0.66	0.659	0.621	0.596	0.559	R
dit07	12	2705	1487	1579	946	1137	593	yr
			0.688	0.611	0.54	0.532	0.51	R
bark11	10	3311	937	873	711	1364	574	yr
			0.745	0.614	0.573	0.573	0.572	R
thresh	11	2682	1475	546	866	1134	594	yr
			0.598	0.574	0.558	0.54	0.536	R
bayes	9	3908	548	1482	1572	652	1167	yr
			0.718	0.692	0.689	0.629	0.616	R

Table 1. Different proposed DO sequences ("Seq.") in question for a) the last glacial period and b) the 42-10 ky interval for the NGRIP record. Number of events considered for the analysis ("#Ev."), the individual estimates of inter-event waiting times ("Estimate") in the exponential distribution which is used for Monte Carlo Simulations. Furthermore, the 5 highest R scored periods for each sequence in decreasing order (P1 to P5). Highlighted in red/blue/green we find the periods in a 10% range around the periods of interest 1470/900/1150. Grey values indicate out of confidence interval R values. The warming events for the different DO-sequences are listed in Table 2. For the Rayleigh Test itself only the $R_{max} = R(P1)$ is considered.

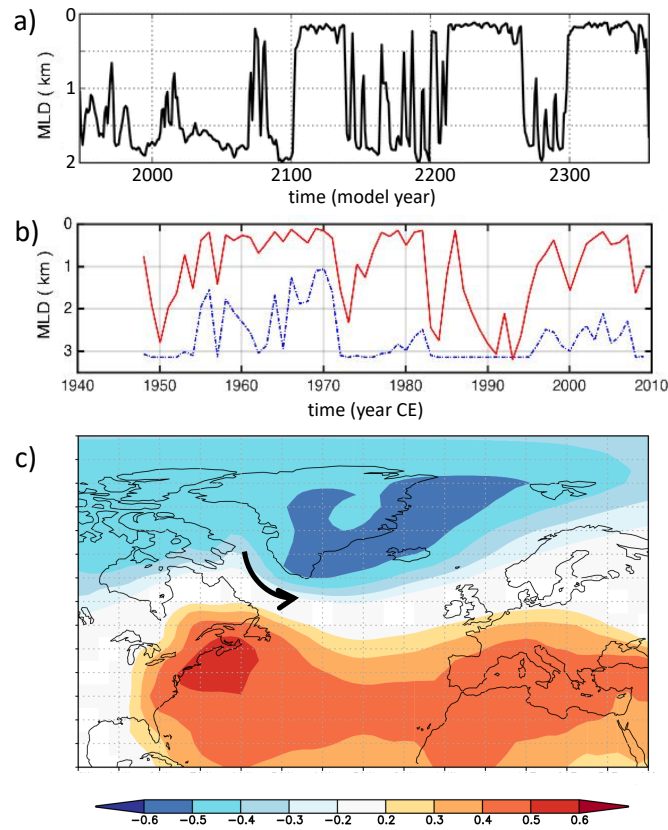


Figure 7. a) Time series of the annual Labrador Sea mixed layer depth (MLD) maximum in a long control integration (Sidorenko et al., 2015) indicating decadal to centennial variability. b) Time series of the March Labrador Sea mixed layer depth for a low (dashed blue) and high (red) resolution in an ocean model forced by atmospheric data (Danek et al., 2019). c) Correlation of the red line in b) with sea level pressure in January using (NCEP reanalysis data, Kalnay et al., 1996). The black arrow shows schematically the surface wind related to deep convection.

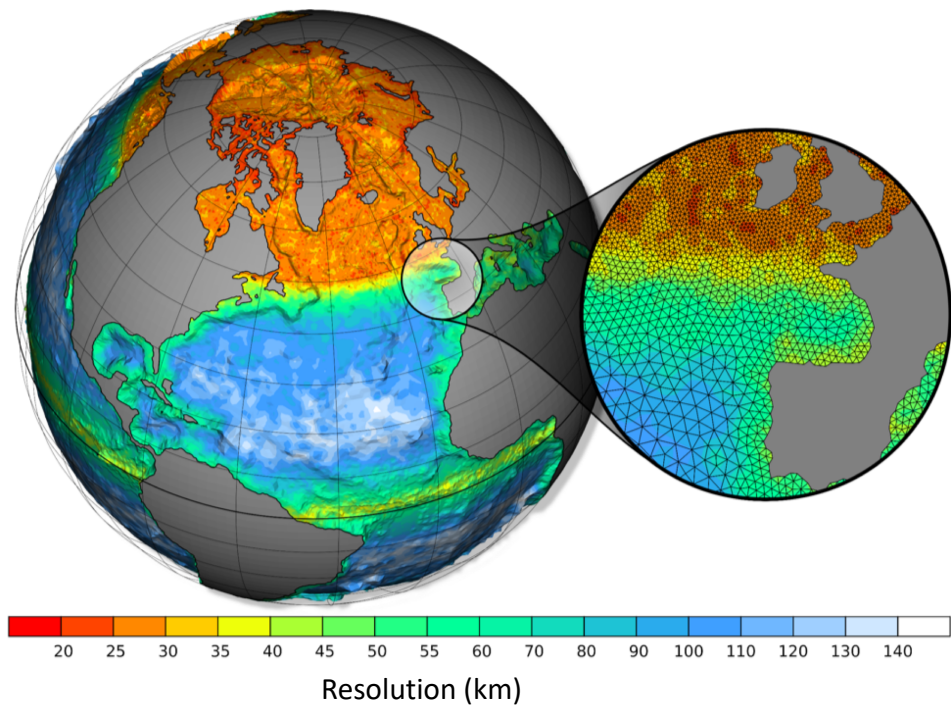


Figure 8. Resolution of the ocean model component in our Earth system model AWI-ESM. Example of an unstructured mesh for our FESOM model with horizontal resolution between 12 and 135 km (CORE-2 mesh, Sidorenko et al., 2011; Sein et al., 2018).

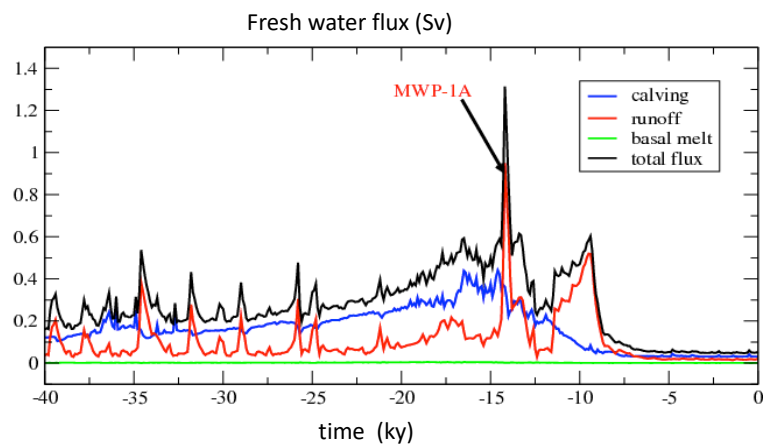


Figure 9. Components of the total freshwater flux (in $\text{Sv} = 10^6 \text{m}^3 \text{s}^{-1}$) originating from the Northern Hemisphere ice sheets at the end of the last ice age based on Zweck and Huybrechts (2005).

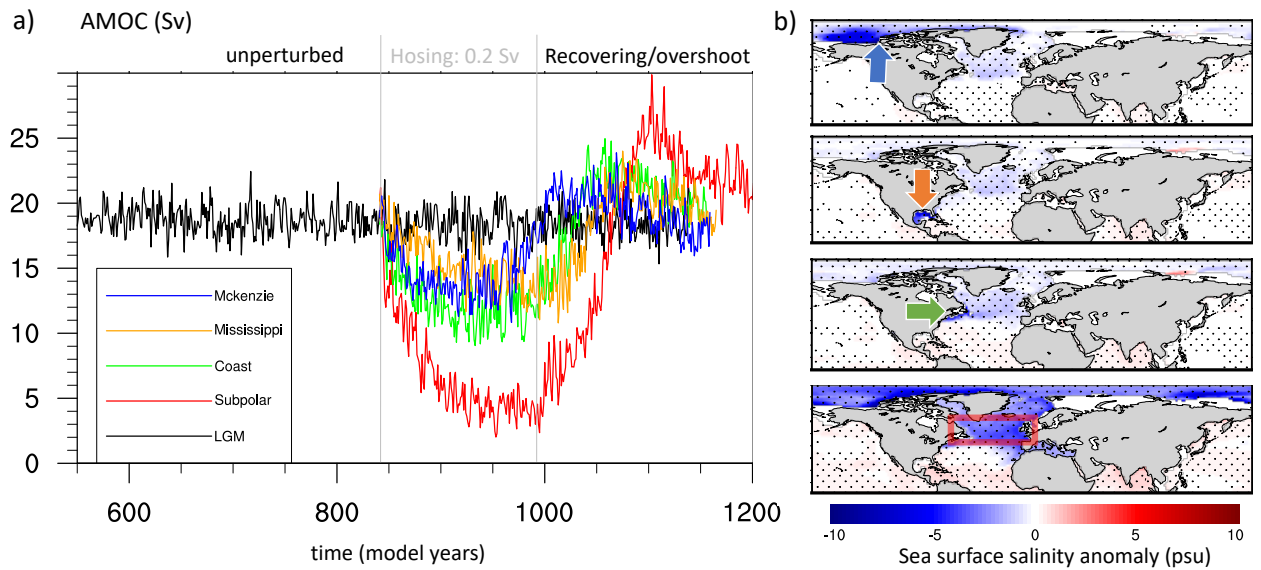


Figure 10. a) AMOC indices of North Atlantic hosing for different hosing areas. Units are Sv. In black line represents the unperturbed LGM experiment. Hosing is for the period 840-990. b) Annual mean sea surface salinity anomaly between LGM and the perturbation experiment LGM with 0.2 Sv for the model years 900-950. Dotted areas are significant on a 95% confidence level (t-test).

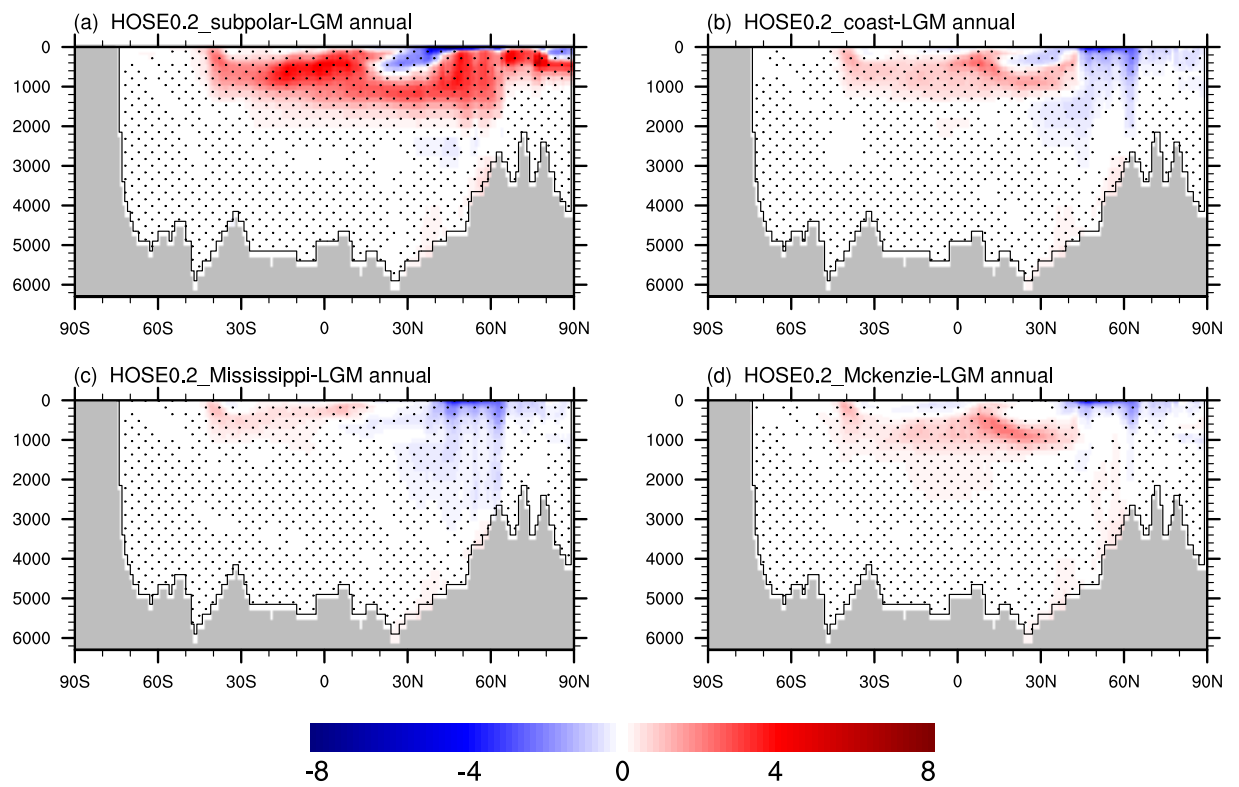


Figure 11. Annual mean Atlantic Ocean temperature anomaly (K) for the different hosing scenarios. a) Subpolar minus LGM, b) Coastal minus LGM, c) Mississippi minus LGM, d) Mackenzie hosing minus LGM. Shown are model averages for the years 900-950. Dotted areas are significant on a 95% confidence level (t-test).

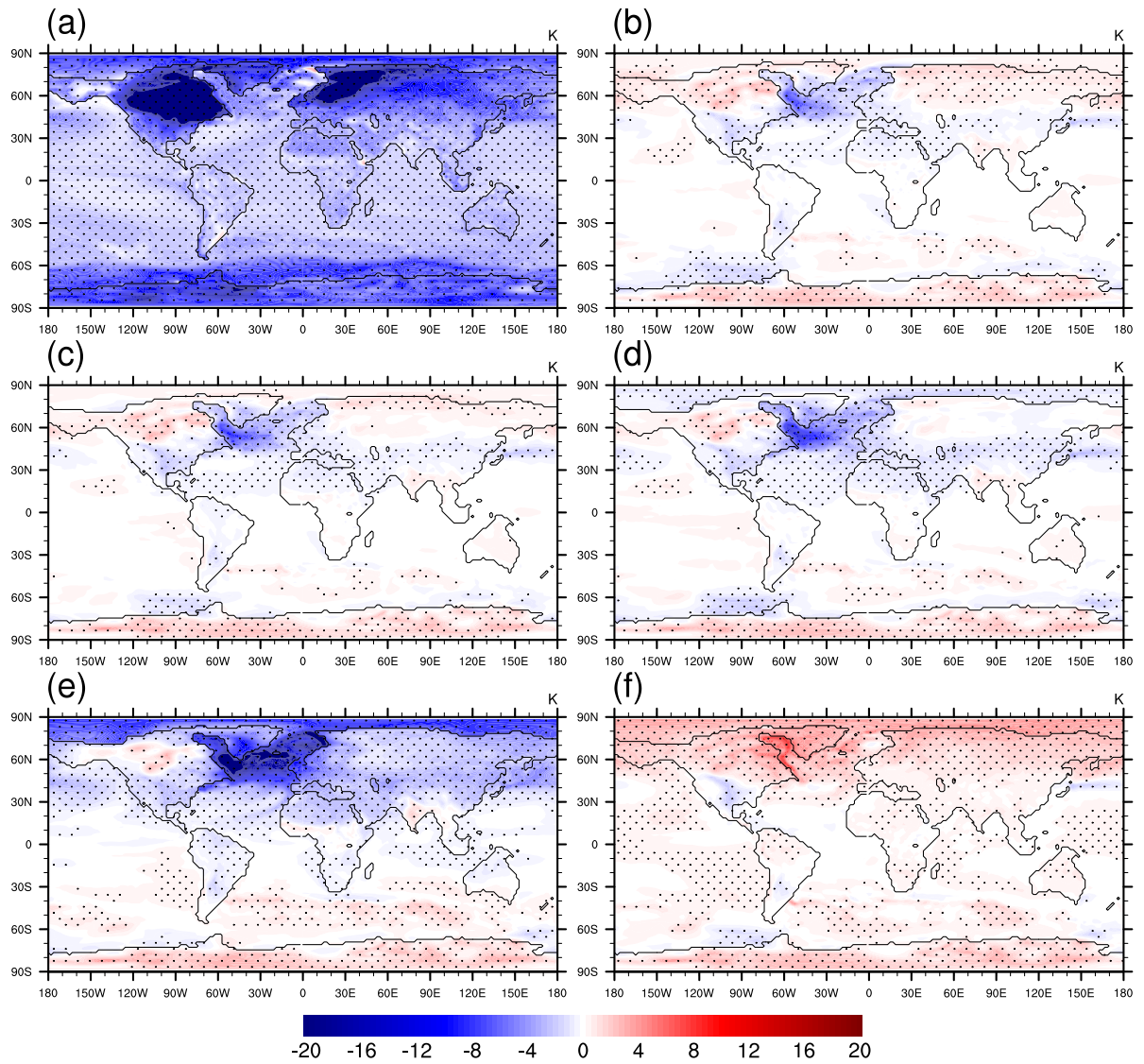


Figure 12. Near-surface air temperature anomaly (K) for a) LGM minus PI, b) Mackenzie hosing minus LGM, c) Mississippi hosing minus LGM, d) Coast minus LGM, e) Subpolar minus LGM, f) Subpolar overshoot minus LGM. Shown are model averages for the years 940-990, and for the overshoot the years 1075-1125. Dotted areas are significant on a 95% confidence level (t-test).

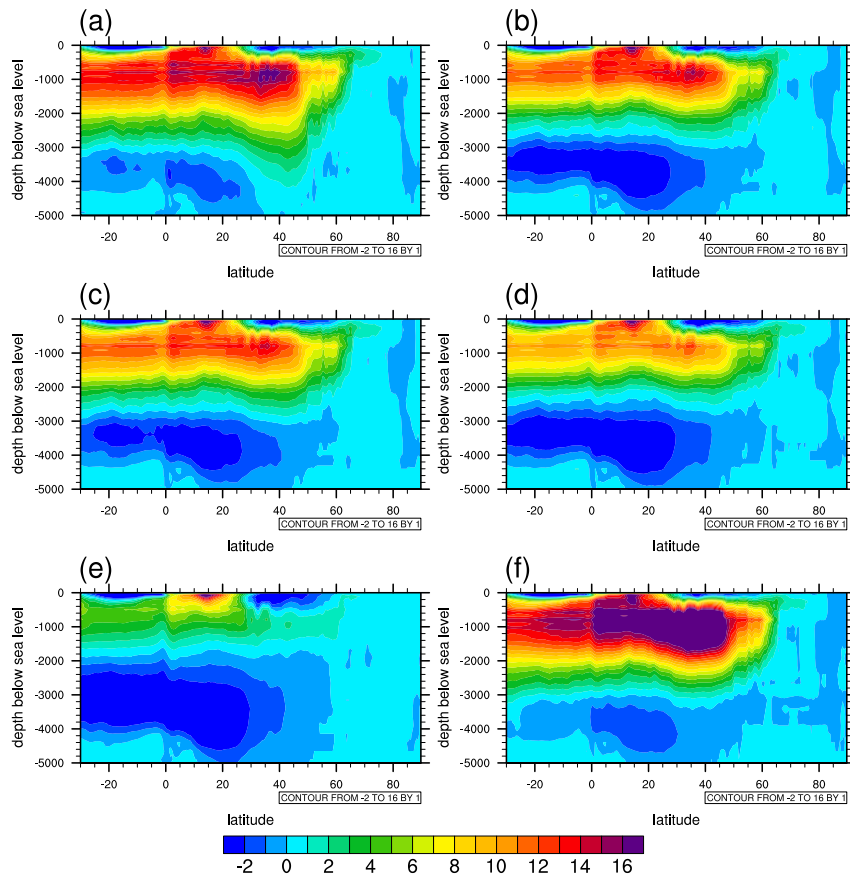


Figure 13. AMOC (Sv) for a) LGM, b) Mackenzie hosing, c) Mississippi hosing, d) Coastal hosing, e) Subpolar hosing, f) Subpolar overshoot. Shown are model averages for the years 940-990, and for the overshoot the years 1075-1125.

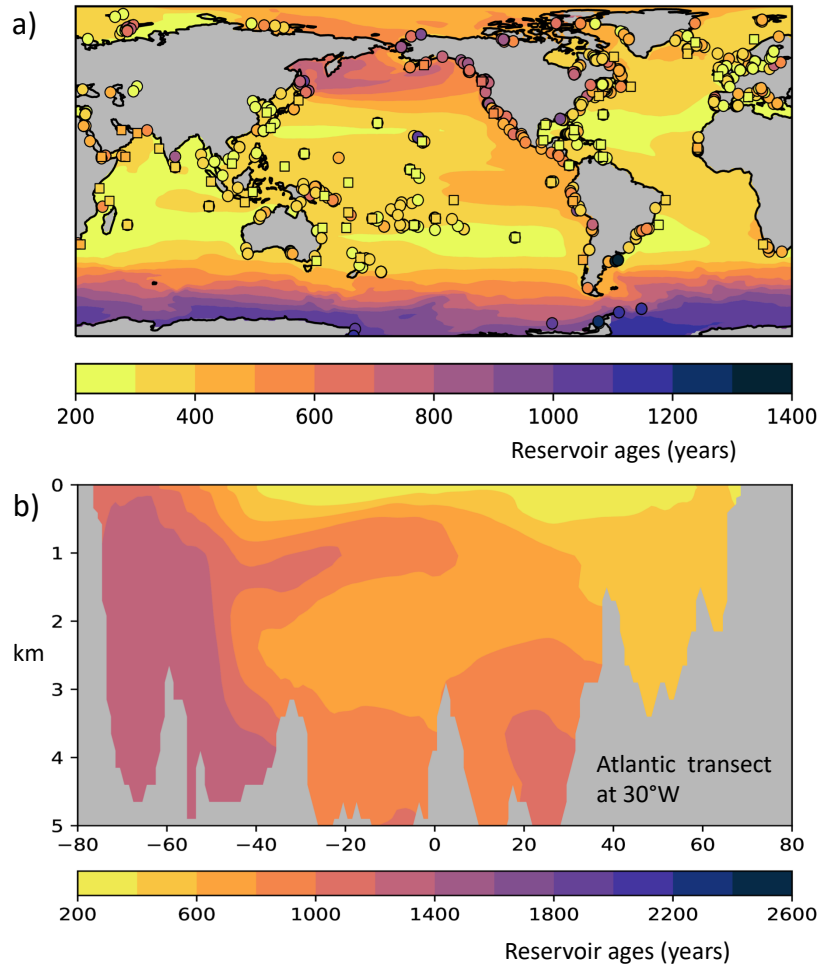


Figure 14. Radiocarbon age in the late Holocene. Filled areas are simulation results using a ^{14}C -equipped version of the global multi-resolution ocean model (Appendix C). a) At the sea surface where dots indicate historical values collected at the ^{14}C CHRONO Marine Reservoir Database (see <http://calib.org/marine/> and further references there). b) Meridional distribution of simulated ^{14}C ages in the Central Atlantic along 30°W .

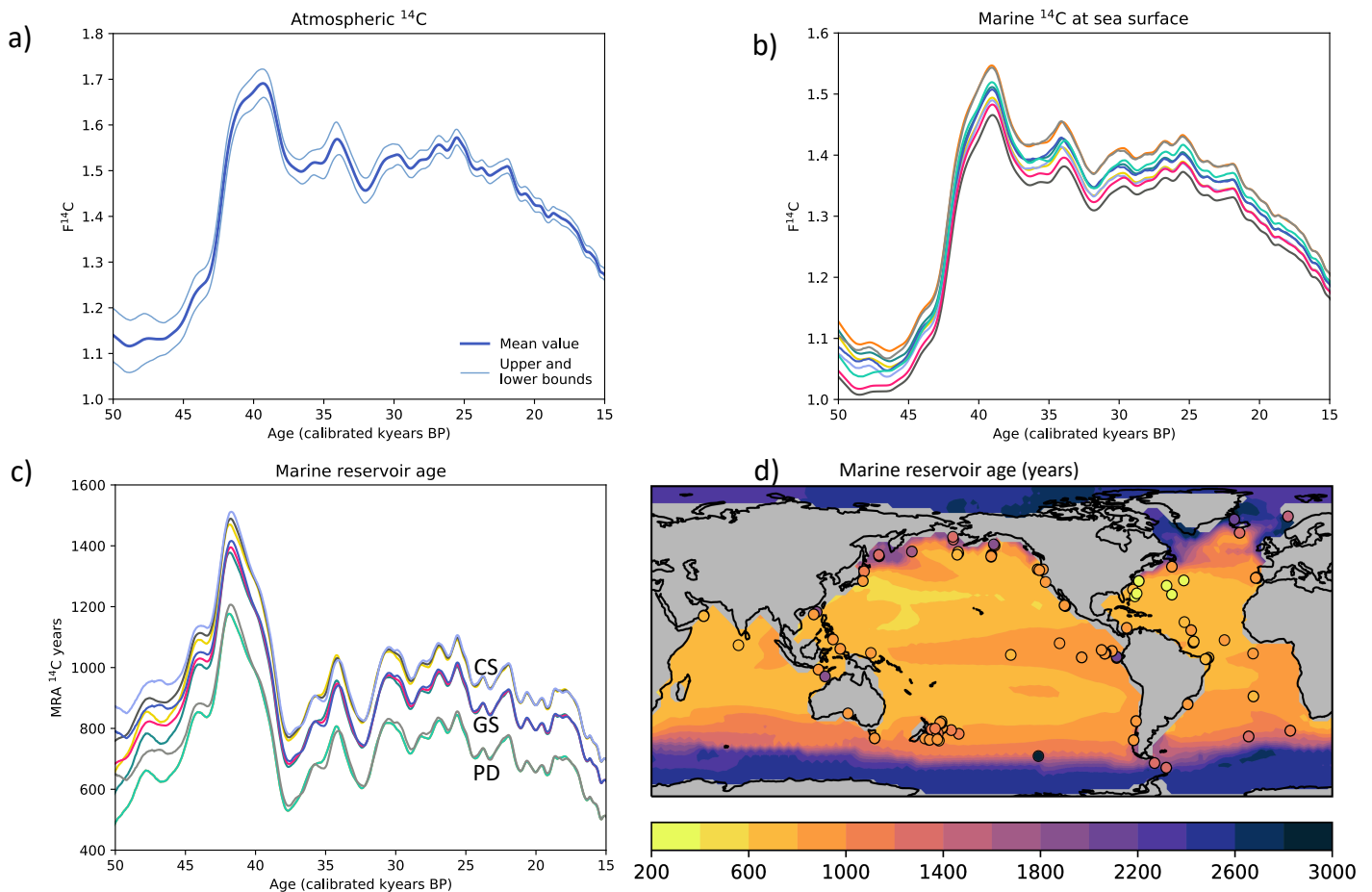


Figure 15. a) Atmospheric radiocarbon concentrations in terms of $F^{14}\text{C}$ (the fractionation-corrected and normalized $^{14}\text{C}/^{12}\text{C}$ ratio) according to the Hulu Cave speleothem record (Southon et al., 2012; Cheng et al., 2018). Upper and lower curves span the uncertainty range (mean values $\pm 2\sigma$). b) Ensemble simulations of marine ^{14}C for the past 50 ky forced with atmospheric $F^{14}\text{C}$ according to Hulu Cave, shown is the period where ^{14}C -dating is not further constrained through continuous tree ring ^{14}C records. c) Corresponding ensemble simulations of marine ^{14}C expressed as ^{14}C age with respect to the contemporaneous atmosphere (the Marine Reservoir Age), values are averaged between $50^\circ\text{N} - 50^\circ\text{S}$. CS, GS, PD specify different ocean states with weak, intermediate, and strong overturning (Butzin et al, 2020). d) Marine reservoir age simulated for the Last Glacial Maximum; shown is the ensemble median of the transient simulations. Filled circles are foraminifera-based marine reservoir ages compiled by Skinner et al. (2017).

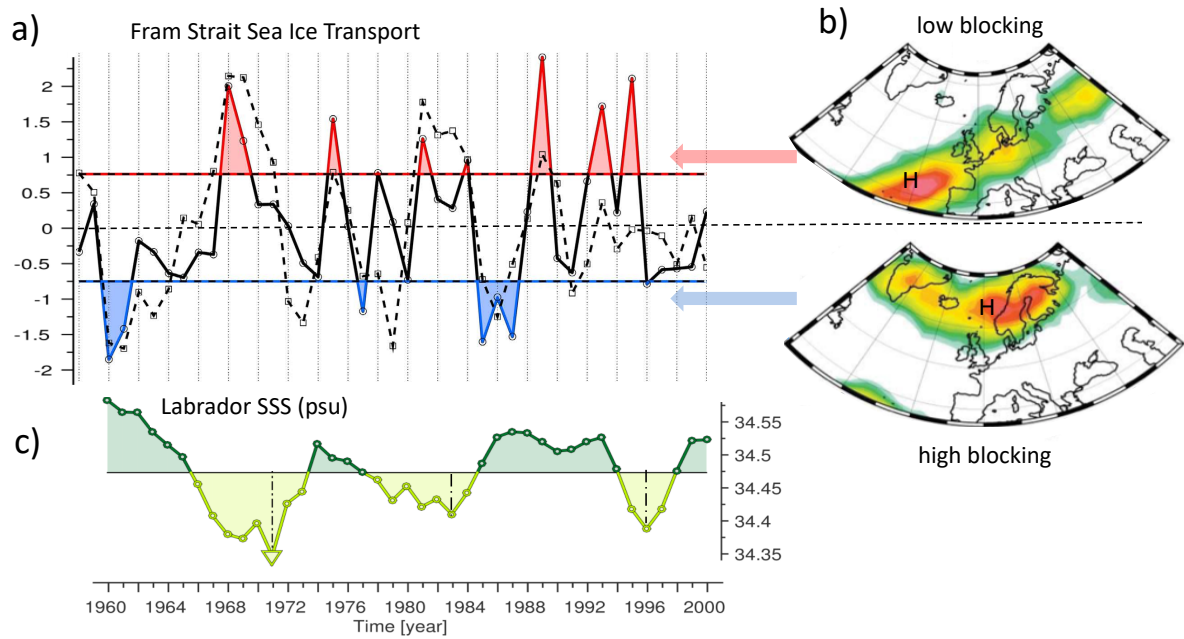


Figure 16. a) Winter Fram Strait sea-ice export time-series (normalized): model simulation (solid line), observations (dashed line). b) Atmospheric blocking frequency composite maps with respect to the simulated Fram Strait sea-ice export: Maps for above and below 75% standard deviation indicating different regimes of blocking in the North Atlantic realm. c) The 3 year-running mean winter sea surface salinity in the central Labrador Sea with three major freshwater events, based on Ionita et al. (2016).

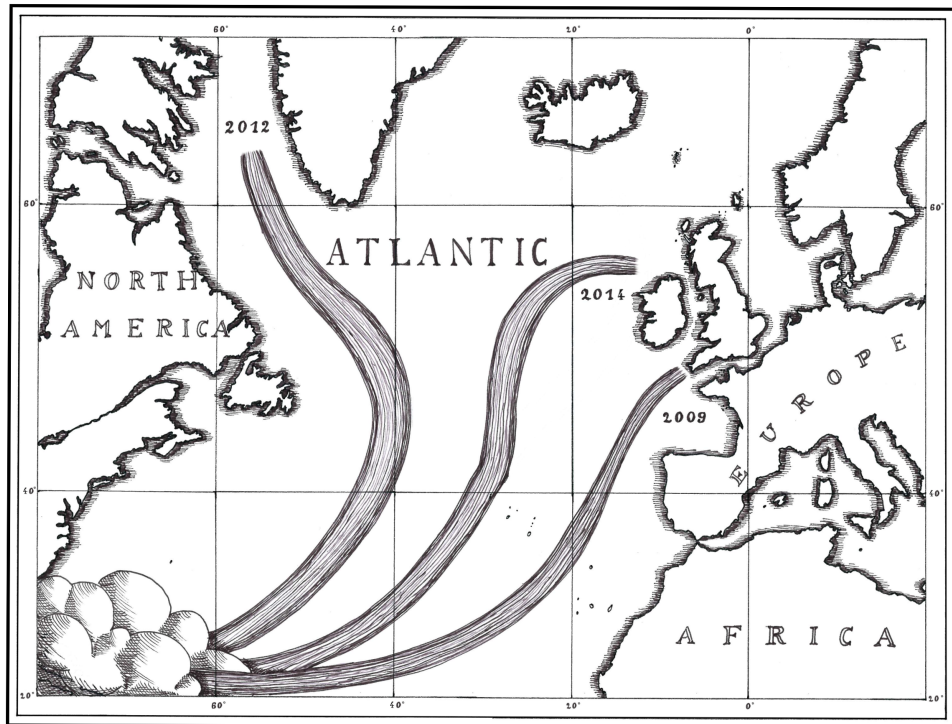
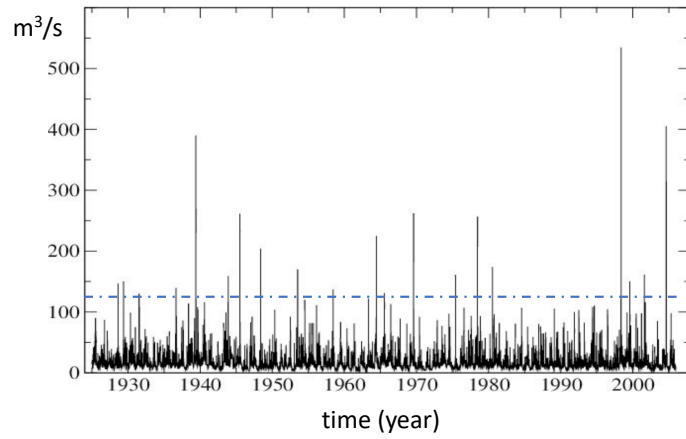


Figure 17. Schematic representation of atmospheric rivers (atmospheric conveyor belt) associated with extreme precipitation events, e.g. in the United Kingdom in 2009.

a)



b)

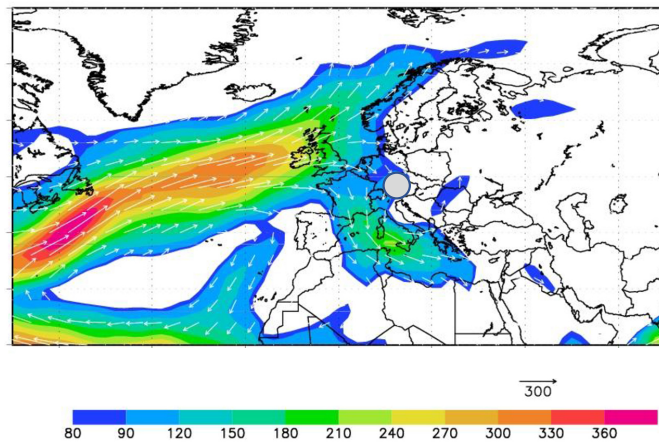


Figure 18. a) Time series of the observed mean daily river Ammer runoff during the period 1926-2006. The dashed line is for daily discharge of $125 \text{ m}^3 \text{ s}^{-1}$. b) Composite map of vertically integrated water vapor transport for flood days (vector) and its magnitude (color). Units: $\text{kg m}^{-1} \text{ s}^{-1}$ for runoff higher than $125 \text{ m}^3 \text{ s}^{-1}$. The geographical Ammer runoff location is inserted as grey dot. Based on Rimbu et al. (2016b).

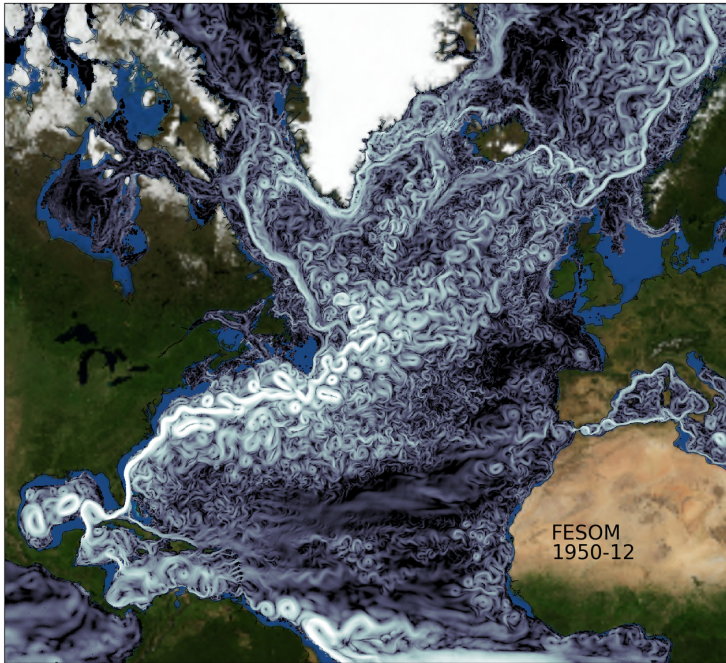


Figure 19. Representation of small scale features of ocean currents in high resolution ocean models: simulated velocity field: simulated velocity field in the North Atlantic at 100 m depth in December 1950 using FESOM with high-resolution locally eddy-resolving mesh (Sein et al., 2018).

878 **Appendix A: COSMOS model**

879 The atmosphere is represented in our model by means of the ECHAM5 atmosphere general circula-
880 tion model (Roeckner et al., 2003). ECHAM5 is based on a spectral dynamical core and resolves in our
881 model setup 19 vertical hybrid sigma-pressure levels. The series of spectral harmonics is curtailed via tri-
882 angular truncation at wave number 31 (approx. 3.75×3.75). Beyond climate system components that
883 are conventionally included in atmosphere general circulation models, ECHAM5 resolves, via its land sur-
884 face component JSBACH (Raddatz et al., 2007), various land surface processes and a representation of
885 vegetation dynamics (Brovkin et al., 2009). Ocean circulation and sea ice dynamics are computed by the
886 MPIOM ocean general circulation model (Marsland et al., 2003) that is employed by us at 40 unequally-
887 spaced levels on a bipolar curvilinear model grid with formal resolution of 3.0×1.8 longitude by lati-
888 tude. The coupled model system ECHAM5/MPIOM is described by Jungclauss et al. (2006). A concise
889 description of the application of the COSMOS for paleoclimate studies is given by Stepanek & Lohmann
890 (2012). The COSMOS version we are using has proven to be a suitable tool for the study of the Earth's
891 past climate, from orbital to tectonic time scales. Examples for applications within the time period of
892 the last 1 My include orbitally forced climate simulations of the current (Wei et al., 2012; Wei & Lohmann,
893 2012; Lohmann et al., 2013a) and previous (Pfeiffer & Lohmann, 2016) interglacials as well as of the Last
894 Glacial Maximum (Zhang et al., 2013).

895 While COSMOS is numerically more efficient than current state-of-the-art climate models, its ac-
896 celerated integration still takes more than 10 real-time calendar months. We note that when exposing
897 the COSMOS to a 100-year sampling of orbital parameters we derive a deep ocean that is less equilibrated
898 to the applied orbital forcing than it actually has been the case during the past. Yet, as stated by Lorenz
899 & Lohmann (2004) who describe and apply this method, even for 100 times accelerated orbital forcing,
900 the necessary conditions for simulating a sensible climate response are fulfilled: First, the applied accel-
901 erated changes in seasonal insolation are small compared to the seasonal cycle. Second, the modes of in-
902 ternal variability in the atmosphere and in the surface ocean are faster than variations in the orbital forc-
903 ing. The standard model code of the Community Earth System Models (COSMOS) version COSMOS-
904 landveg r2413 (2009) is available upon request from the Max Planck Institute for Meteorology in Ham-
905 burg (<https://www.mpimet.mpg.de>). The temperature data of COSMOS related to Figures 4 and 5 are
906 available through PANGAEA (Stepanek and Lohmann, 2020).

Appendix B: Defining DO - events with the Bayesian Change Point Analysis

For the identification of DO-event dating of the last glacial is based on

- (1) a time sequence based on a threshold criterion defined by Ditlevsen et al. (2005), followed by
- (2) a Bayesian approach which allows us to verify choices of events with uncertainties.

We will consider exact dating sequences used in previous studies. As DO-events are similar in shape and characterized by the abrupt warming we will reduce the single events to single time points of warming onset. Criterion (1) goes along with the assumption that DO events represent repeated oscillations between interstadial and stadial conditions with a threshold crossing process involved (Alley et al., 2003; Ditlevsen et al., 2007; Steffensen et al., 2008). We consider a threshold criterion for the definition of DO events. In accordance with Ditlevsen et al. (2005) we define a time point an event at the first upcrossing of a higher threshold (initiation of an interstadial state), if it is consequently followed by an up-crossing of a lower threshold (relaxation process to stadial mode). A trend for the scaled NGRIP ice core is estimated by single spectral analysis (SSA) and subtracted, the higher threshold is set to 1 and the lower threshold is chosen to be -0.9 anomaly from the estimated trend. The asymmetry is justified by the difference in state duration, stadial modes persist longer (Ditlevsen et al., 2005).

The Bayesian Change Point Analysis provides us with uncertainty estimates of the number of change points and allows us to input a priori knowledge (Ruggieri, 2013). It splits the problem of possible placements of change points into sub problems of smaller data chunks and uses a recursive dynamic programming for reduction in computational costs. For this study we only consider those time points as DO-events that result a certainty of location greater than 0.5. For the NGRIP record this results in the choice of 22 events (Fig. 6, Table 2).

DO – Time Sequences					
Event	Ras14	Dit07	Bark11	Thresh	Bayes
0	11703	11700	11500	-	10200
A	-	13130	-	-	-
1	14692	14680	14300	14620	14680
B	-	-	-	16520	-
C	-	-	-	20740	-
2	23340	23340	-	23340	-
3	27780	27780	27380	27780	27780
4	28900	28900	29260	28880	28880
5	32500	32500	32020	32500	32500
6	33740	33740	33340	33720	33740
7	35480	35480	35020	35480	35500
8	38220	38220	37780	38200	38220
9	40160	40160	38680	-	-
10	41460	41460	41300	41440	41460
11	43340	-	43100	43127.1	43386.3
12	46860	-	46420	47053.5	47053.5
13	49280	-	-	-	-
14	54220	-	53300	54434.7	54754.1
15	55800	-	-	55503.2	-
16	58280	-	57080	-	-
17	59440	-	58580	59508.3	59876.7
18	64100	-	63840	64477.5	64477.5
19	72340	-	70680	72689.1	72743.9
20	76440	-	74500	76821.5	76821.5
21	85060	-	83180	85074.8	85155.4
22	90040	-	89920	90092.3	90449.5
23	104520	-	101520	104251	104412
24	108280	-	106540	108331	108435
25	115370	-	-	115111	115163
Total number	26	12	22	24	21

Table 2. Timing of the DO events for the last glacial found in the NGRIP Ice Core of Greenland (GICC05 time scale). The timings are taken from publications Rasmussen et al. (2014), Ditlevsen et al. (2007) and Barker et al. (2011). Additionally obtained from thresholding and Bayesian estimation are given. For the total number of events we do not consider the start of the Holocene (event 0).

928 Appendix C: The multi-scale ocean circulation-radiocarbon model

929 The AWI Earth System Model (AWI-ESM) is an extension of the AWI-CM (Sidorenko et al., 2015,
 930 2019) for earth system modelling (<https://fesom.de/models/awi-esm/>). The atmospheric module is rep-
 931 resented by the general circulation model ECHAM6 (Stevens et al., 2013) (here in T63 horizontal res-
 932 olution, $\sim 180km$) including a Land-Surface Model (JSBACH) which is based on a tiling of the land sur-
 933 face and includes dynamic vegetation with 12 plant functional types and two types of bare surface (Rad-
 934 datz et al., 2007). The ocean and sea ice component is the Finite Element Sea Ice-Ocean Model (FESOM)
 935 (Danilov et al., 2004; Timmermann et al., 2009; Sidorenko et al., 2011; Rackow et al., 2018a, 2019) which
 936 is discretized on a triangular grid with a continuous conforming representation of model variables. The
 937 mesh nodes are vertically aligned to avoid difficulties in resolving the hydrostatic balance. The model uses
 938 variable resolution, which can reach 20 km in the Arctic and along coastlines. A no-slip boundary con-
 939 dition along the coast is implemented in the model (Fig. 8). Surface stress and buoyancy fluxes are de-
 940 rived from the ice-ocean coupling. FESOM model has been validated in Timmermann et al. (2009) and
 941 Scholz et al. (2013), and the coupled AWI climate model in more recent studies (Sidorenko et al., 2015;
 942 Rackow et al., 2019). The source code of the climate model is available from the AWI based svn repos-
 943 itory (<https://swrepo1.awi.de/projects/awi-cm/>).

944 The radiocarbon-climate model is based on FESOM2 (Danilov et al., 2017). This model has been
 945 used as the ocean-sea ice component in the AWI climate model (Sidorenko et al., 2019). The current ver-
 946 sion of FESOM2 is available from the public GitHub repository at <https://github.com/FESOM/fesom2>
 947 under the GNU General Public License (GPLv2.0). Radiocarbon is treated as $F^{14}C$ following Toggweiler
 948 et al. (1989), with an air-sea gas exchange formulation accounting for glacial climatological boundary con-
 949 ditions. Marine biological processes are neglected because these effects play a minor role for $F^{14}C$ com-
 950 pared to the changes induced by circulation and radioactive decay (Fiadeiro, 1982). Radiocarbon data
 951 are frequently quoted in the form of ages, via $^{14}C_{age} = t_{1/2} \cdot \ln(^{14}F_a / ^{14}F_o) / \ln 2$, where $t_{1/2}$ is the half-
 952 life of ^{14}C , and $^{14}F_a$ and $^{14}F_o$ are the normalized and fractionation corrected $^{14}C/^{12}C$ ratios in atmo-
 953 sphere and ocean. High radiocarbon concentrations in water translate into low radiocarbon ages and vice
 954 versa. Note that "conventional" radiocarbon age values are based on the "Libby half-life" of 5568 years
 955 (Stuiver & Polach 1977), while the most recent estimate of the true half-life is 5700 years (Audi et al.,
 956 2003, Bé et al., 2013).

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